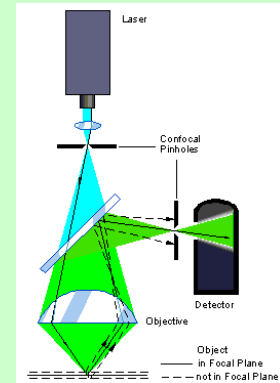
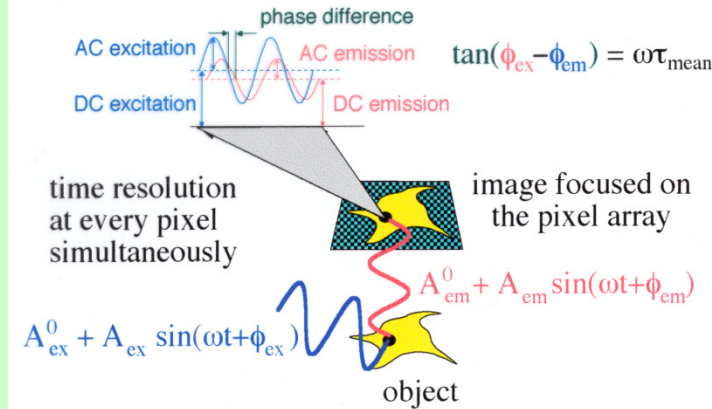
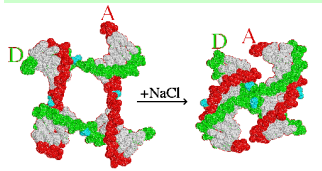


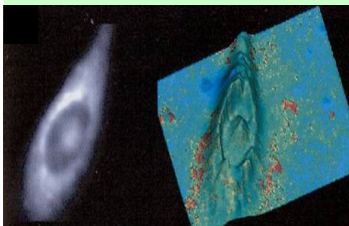
Organisms



The marriage of morphology and lifetimes:
Excavating the tortuous landscape of FLIM



Molecules



Cells

τ_1

τ_2

WITH A PULLY SYSTEM, YOU CAN HOLD BOONE, WHO WEIGHS MUCH MORE THAN YOU!

OF COURSE YOU ALWAYS KEEP TWO HANDS ON THE ROPE!

τ_1

τ_2



Medical Imaging

**Measuring lifetimes has
been around
for a long time**

And

The theory behind it too.



Si la loi énoncée plus haut est exacte, on doit avoir d'après la formule (2)

$$(3) \quad \log i = \log i_0 - at, \quad \log i' = \log i_0 - at', \text{ etc...}$$

d'où

$$\frac{\log i' - \log i}{t - t'} = a, \quad \frac{\log i'' - \log i'}{t' - t''} = a, \dots,$$

c'est-à-dire que les différences entre les logarithmes des intensités lumineuses doivent être proportionnelles aux différences des temps, et que leur rapport doit donner précisément le coefficient a .

$$\frac{di}{dt} = -ai,$$

$$(2) \quad i_t = i_0 e^{-at},$$

$$Q = \int_0^{\infty} i_0 e^{-at} dt = \frac{i_0}{a},$$

$$I = i_0 e^{-at} + y_0 e^{-bt},$$

LA LUMIÈRE

SES CAUSES ET SES EFFETS

PAR

(Alexandre)

M. EDMOND BECQUEREL

DE L'ACADÉMIE DES SCIENCES

DE L'INSTITUT DE FRANCE

PROFESSEUR AU CONSERVATOIRE IMPÉRIAL DES ARTS ET MÉTIERS, ETC., ETC.

TOME PREMIER

PARIS

LIBRAIRIE DE FIRMIN DIDOT FRÈRES, FILS ET C^{IE}

IMPRIMEURS DE L'INSTITUT, RUE JACOB, 56

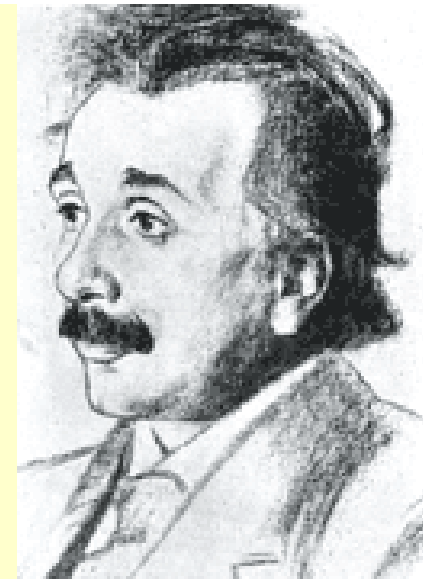
1867

$$g_1 B_{12} = g_2 B_{21}$$

$$\frac{1}{\tau} \Rightarrow$$

$$A_{21} = 16\pi^2 \hbar (\omega/2\pi c)^3 B_{21}$$

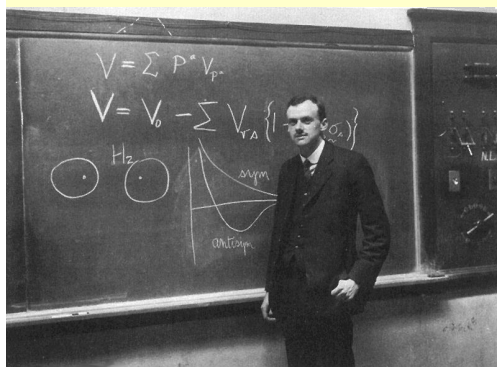
$$\frac{A_m^n}{B_m^n} = \alpha \nu^3$$



ON THE QUANTUM THEORY OF RADIATION

¹ A. Einstein. Physik Z. 18. 121 (1917)

IN a classic paper, Einstein¹ described relations connecting the rates of spontaneous emission, stimulated emission, and absorption of radiation by an atomic system in free space having two sharp energy levels.



$$\frac{1}{\tau} \Rightarrow \lambda_{if} = \frac{2\pi}{\hbar} |M_{if}|^2 \rho_f$$

Transition probability
Matrix element for the interaction
Density of final states

Fermi's Golden Rule

natural radiative lifetime

The Quantum Theory of the Emission and Absorption of Radiation P. A. M. Dirac

Proceedings of the Royal Society of London. Series A, Containing Papers of a Mathematical and Physical Character, Vol. 114, No. 767 (Mar. 1, 1927), 243-265.

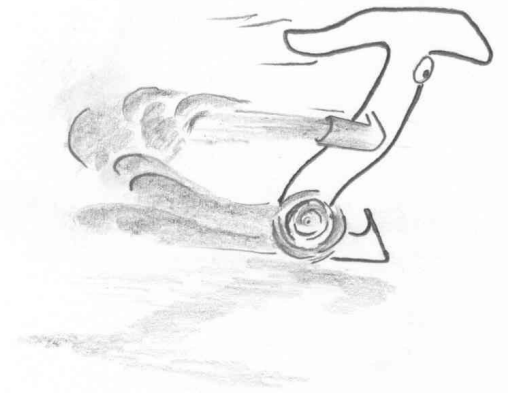
The probability per unit time of a transition to a state for which each γ_k lies between γ_k' and $\gamma_k' + d\gamma_k'$ is thus (apart from the normalising factor)

$$\frac{1}{\tau} \Rightarrow$$

$$2\pi |\alpha^0|^2 / \hbar \cdot |v(W^0, \gamma'; W^0, \gamma^0)|^2 J(W^0, \gamma') d\gamma_1' \cdot d\gamma_2' \dots d\gamma_{u-1}', \quad (24)$$

which is proportional to the square of the matrix element associated with that transition of the perturbing energy.

How does one measure
nanosecond lifetimes?



$$F(t)_{meas} = \int_0^t E(t') F_{\delta}(t-t') dt'$$

So called:

Time-domain:

$$F_{\delta}(t-t')_{meas} = \sum_i F_{\delta,i}(t-t') = \sum_i F_{0,i} \exp(-(t-t')/\tau_i)$$

Frequency domain:

Excitation repetitive pulse; e.g. $\rightarrow \propto \cos(\omega t)$

$$F(t)_{meas} = \left[\sum_i F_{0,i} \tau_i + \sum_i \frac{F_{0,i} \tau_i}{1+j\omega\tau_i} e^{j\omega t} \right] = \left[\sum_i F_{0,i} \tau_i + e^{j\omega t} \sum_i \frac{F_{0,i} \tau_i}{\sqrt{1+(\omega\tau_i)^2}} e^{-j \tan^{-1} \omega\tau_i} \right]$$

$$\frac{F(t)_{meas}}{F_{meas,ss}} = 1 + \sum_i \frac{\alpha_i}{1+j\omega\tau_i} e^{j\omega t} = 1 + e^{j\omega t} \sum_i \alpha_i M_i \left[\cos(\phi_{i,\omega}) + j \sin(\phi_{i,\omega}) \right]$$

$$M = \frac{1}{\sqrt{1+(\omega\tau_M)^2}} \quad \phi = \tan^{-1}(\omega\tau_{\phi})$$

What can affect the lifetime?

**Let's see how
the old masters measured it**

Die Abklingungszeiten der Fluoreszenz von Farbstofflösungen.

Von E. Gaviola in Berlin.

(Eingegangen am 10. Dezember 1925.)

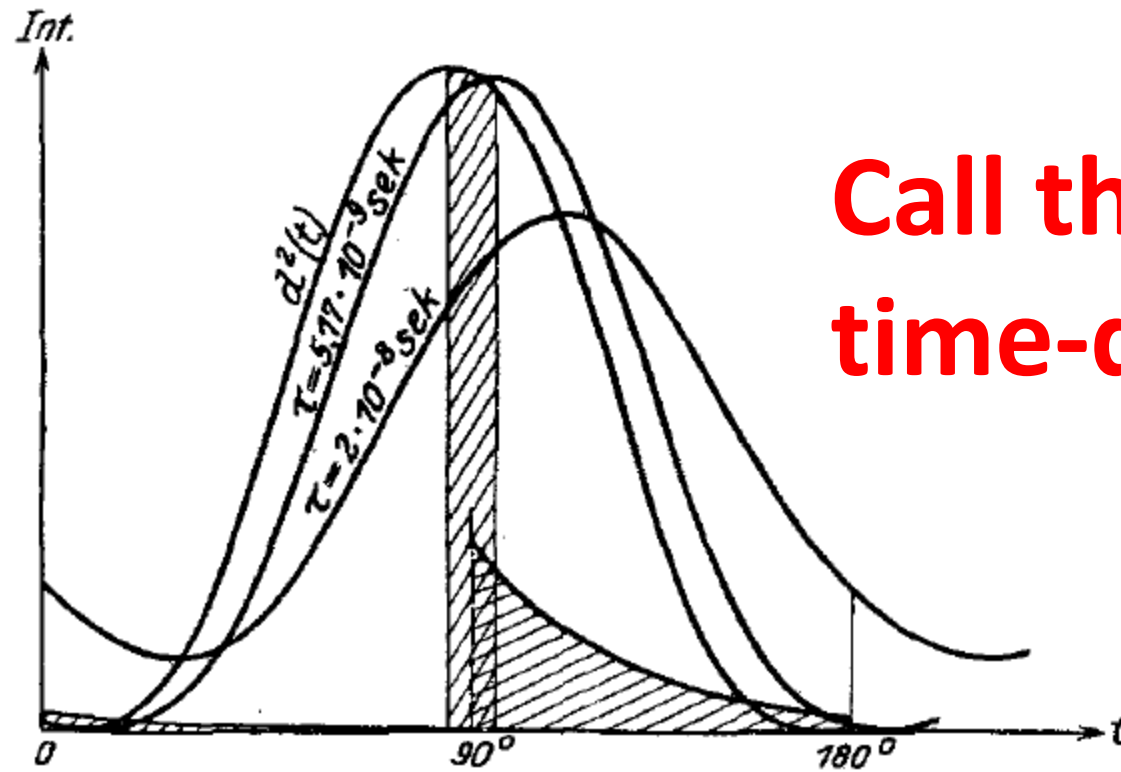


Fig. 6.

Call this the
time-domain

The signal is convoluted
with the excitation pulse

$$K(t^*) = \int_{-\infty}^{t^*} \frac{d^2(t)}{\tau} e^{-\frac{t^*-t}{\tau}} \cdot dt$$



Die Abklingungszeiten der Fluoreszenz von Farbstofflösungen.

Von E. Gaviola in Berlin.

1925

It was difficult to measure, but look at the results!

Farbstoff	Abklingungszeiten		
	in Wasser Sekunden	in Glycerin Sekunden	in Meth.-Alkohol Sekunden
Uranin	$4,5 \cdot 10^{-9}$	$4,4 \cdot 10^{-9}$	—
Fluorescein	—	—	$5,0 \cdot 10^{-9}$
Rhodamin B	$2,0 \cdot 10^{-9}$	$4,2 \cdot 10^{-9}$	—
Rhodulin Orange	2,7	4,3	—
Erythrosin	1,8	2,4	$2,6 \cdot 10^{-9}$
Tetraiodfluor. Na	1,0	2,0	2,2
Eosin 5 B	1,9	—	3,4
Uranylsulfat	—	—	1,3
Uranylsulfat in Schwefelsäure	—	—	1,9
Chinizarin in Pentan	—	—	2,9
Uranglas	—	—	> 15,0
Rubinkristall	—	—	> 15,0

Den mittleren Fehler der oben angegebenen Zahlen schätze ich zu etwa $\pm 0,5 \cdot 10^{-9}$ sec. Er kann unter Umständen viel kleiner sein. Die

**Eine allgemeine Theorie
der zur Messung sehr kurzer Leuchtdauern dienenden
Versuchsanordnungen (Fluorometer).**

Von **F. Duschinsky** in Berlin.

Mit 3 Abbildungen. (Eingegangen am 10. Januar 1933.)

¹⁾ F. Duschinsky, ZS. f. Phys. **81**, 23, 1933.

"Fourier's theorem is not only one of the most beautiful results of modern analysis, but it is said to furnish an indispensable instrument in the treatment of nearly every recondite question in modern physics...

Fourier is a mathematical poem." Lord Kelvin



([March 21, 1768](#) - [May 16, 1830](#))

$$f(t) = \frac{a_0}{2} + \sum_{n=1}^{\infty} \left[a_n \cos \frac{2n\pi t}{T} + b_n \sin \frac{2n\pi t}{T} \right]$$

$$a_0 = \frac{2}{T} \int_0^T f(t) dt$$

$$a_n = \frac{2}{T} \int_0^T f(t) \cos\left(\frac{2n\pi t}{T}\right) dt$$

$$b_n = \frac{2}{T} \int_0^T f(t) \sin\left(\frac{2n\pi t}{T}\right) dt$$

Der zeitliche Intensitätsverlauf von intermittierend angeregter Resonanzstrahlung.

Von F. Duschinsky in Berlin.

(Eingegangen am 10. Januar 1933.)

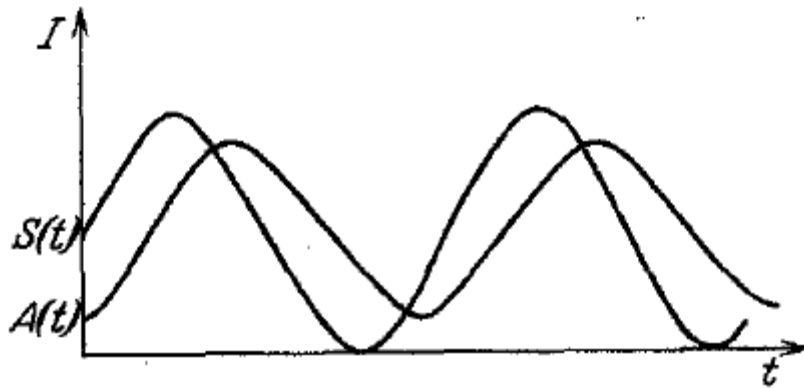


Fig. 1.

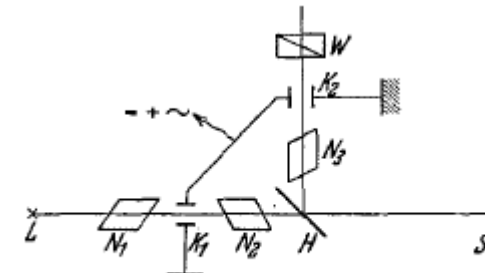


Fig. 1.

Call this the

frequency-domain

Every frequency component
is analyzed separately

$$\left. \begin{aligned} S(t) &= X_0 + \sum_1^{\infty} X_m \cos(m\Omega t - \eta_m), \\ A(t) &= X_0 + \sum_1^{\infty} \frac{X_m}{\sqrt{1 + (m\Omega\tau)^2}} \cdot \cos(m\Omega t - \eta_m - \operatorname{arctg} m\Omega\tau). \end{aligned} \right\} (46')$$

Eine allgemeine Theorie der zur Messung sehr kurzer Leuchtdauern dienenden Versuchsarrordnungen (Fluorometer).

Von **F. Duschinsky** in Berlin.

Mit 3 Abbildungen. (Eingegangen am 10. Januar 1933.)

¹⁾ F. Duschinsky, ZS. f. Phys. **81**, 23, 1933.

$$L(t') = \int_0^{t'} E(t' - t) \Phi(t) dt$$

Ist die Erregungsintensität $E(t)$ periodisch (mit der Frequenz ω' moduliert), so kann sie als FOURIER-Reihe dargestellt werden:

$$E(t) = \sum_{\mu=0}^{\infty} E_{\mu} \cos(\mu \omega' t + e_{\mu}).$$



Herrn Prof. Pringsheim meinen besonderen Dank ausdrücken für alles, was ich von ihm an wissenschaftlicher Anregung erhalten habe. Ebenso bin ich Herrn Prof. Schrödinger für freundlichen Rat und wertvolle Kritik an dieser Arbeit zu großem Dank verpflichtet.



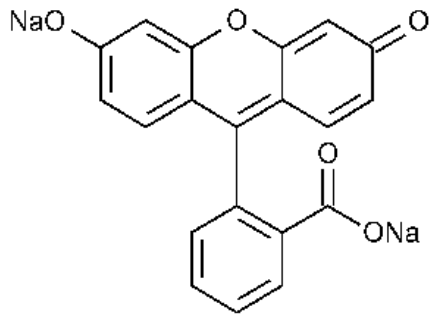
Fig. 1. Schema der verbesserten Fluorometeranordnung. L Lichtquelle; N_1, N_2, N_3, N_4 Nicols, K_1, K_2 Kerr-Zellen; K Babinet-Soleil-Kompensator; S Szivessi-Platte; Z Spiegel; T Trog.

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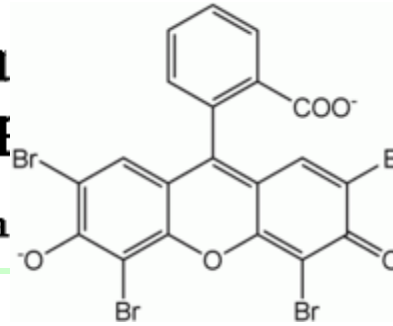
72

The lifetimes do
funny things
depending on
the molecular species
and
the environment



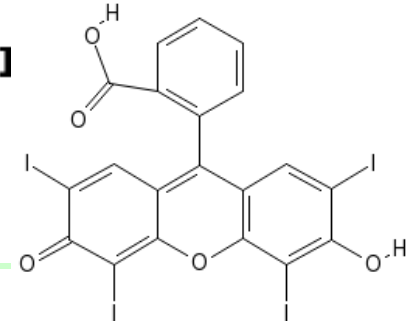
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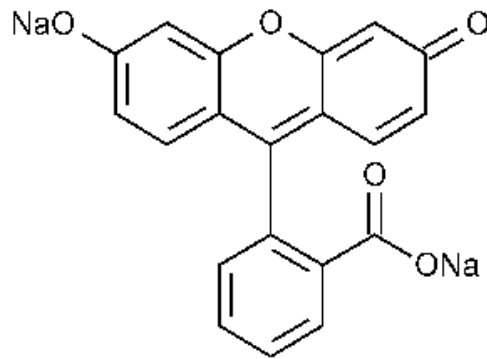


5

It was difficult to measure, but look at the results!
Fluorescein Eosin Y Erythroscine

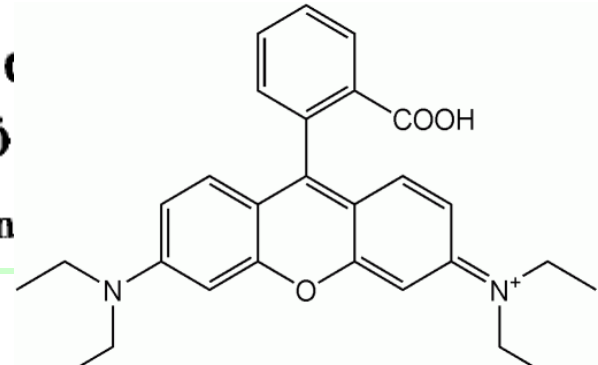
Farbstoff	Abklingungszeiten		
	in Wasser Sekunden	in Glycerin Sekunden	in Meth.-Alkohol Sekunden
Uranin	$4,5 \cdot 10^{-9}$	$4,4 \cdot 10^{-9}$	$5,0 \cdot 10^{-9}$
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Tetraiodfluor . Na	1,0	2,0	2,2
Eosin 5 B	1,9	—	3,4
Uranylsulfat	—	—	1,3
Uranylsulfat in Schwefelsäure	—	—	1,9
Chinizarin in Pentan	—	—	2,9
Uranglas	—	—	> 15,0
Rubinkristall	—	—	> 15,0

Den mittleren Fehler der oben angegebenen Zahlen schätze ich zu etwa $\pm 0,5 \cdot 10^{-9}$ sec. Er kann unter Umständen viel kleiner sein. Die



Abklingzeiten (Farbstofflösungen)

Prof. Gaviola in



325

It was a direct measurement, but look at the results!
Fluorescein **Rhodamine B**

Farbstoff	Abklingungszeiten		
	in Wasser Sekunden	in Glycerin Sekunden	in Meth.-Alkohol Sekunden
Uranin	$4,5 \cdot 10^{-9}$	$4,4 \cdot 10^{-9}$	—
Fluorescein	—	—	$5,0 \cdot 10^{-9}$
Rhodamin B	$2,0 \cdot 10^{-9}$	$4,2 \cdot 10^{-9}$	—
Rhodulin Orange	2,7	4,3	—
Erythrosin	1,8	2,4	$2,6 \cdot 10^{-9}$
Tetraiodfluor. Na	1,0	2,0	2,2
Eosin 5 B	1,9	—	3,4
Uranylsulfat	—	—	1,3
Uranylsulfat in Schwefelsäure	—	—	1,9
Chinizarin in Pentan	—	—	2,9
Uranglas	—	—	> 15,0
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Den mittleren Fehler der oben angegebenen Zahlen schätze ich zu etwa $\pm 0,5 \cdot 10^{-9}$ sec. Er kann unter Umständen viel kleiner sein. Die

**The lifetimes depend on
the dynamics of the molecules
and the dynamics
of the environment**

Die Lebensdauer der angeregten Moleküle in den wässrigen fluoreszierenden Lösungen.

Von S. I. Wawilow in Moskau.

Zeitschrift für Physik A Hadrons and Nuclei
Volume 53, Numbers 9-10 / September, 1929 665-674

Measuring tau using translational diffusion: dynamic quenching

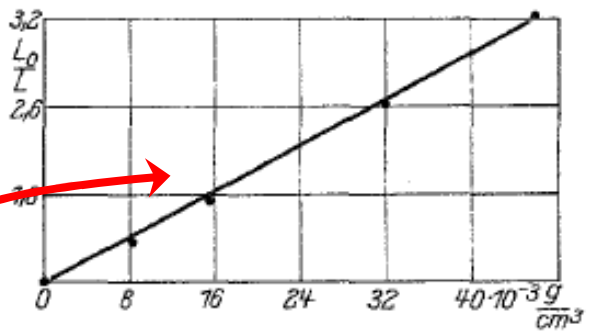
Tabelle 5. Rhodamin B + KJ.

$c_{KJ} \dots$	0	$8,1 \cdot 10^{-3} \text{ g/cm}^3$	15,6	32	46
$L_0/L \dots$	1,00	1,36	1,75	2,64	3,24

$$\bar{t} = \frac{1}{8 \pi D \sigma N c}$$

$$D_1 = \frac{kT}{6 \pi \sigma_1 \eta}, \quad D_2 = \frac{kT}{6 \pi \sigma_2 \eta}$$

$$\frac{L_0}{L} = 1 + \frac{\tau N c k T}{3 \eta} \cdot \alpha,$$



anstatt σ muß $\frac{\sigma_1 + \sigma_2}{2}$

anstatt D $\left(\frac{D_1 + D_2}{2}\right)$

Auch hierbei ist (4) erfüllt, und aus der Neigung der Geraden in der Fig. 5 wird

$$\tau = 2,2 \cdot 10^{-9} \text{ sec}$$

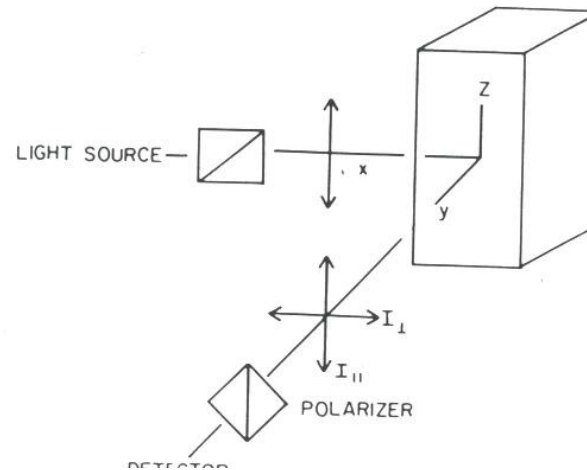
gefunden. Die direkten Messungen von Gaviola für Rhodamin im Wasser ergaben*

$$\tau = 2,0 \cdot 10^{-9} \text{ sec.} \quad * \text{ ZS. f. Phys. 42, 861, 1927.}$$



Francis Henri Perrin
1901-92

Using fast rotational motions to measure the fluorescence lifetime



$$p = \frac{I_{par} - I_{perp}}{I_{par} + I_{perp}}$$

$$r = \frac{I_{par} - I_{perp}}{I_{par} + 2I_{perp}}$$

§ 3. Die Polarisation der Fluoreszenz gibt ein anderes Mittel zur Bestimmung von τ^* . Nach der Formel von F. Perrin ist

$$\frac{1}{p} = \frac{1}{p_0} + \left(\frac{1}{p_0} - \frac{1}{3} \right) \tau \frac{RT}{v\eta}, \quad (7)$$

wo p der Polarisationsgrad bei gegebenen Bedingungen ist, p_0 der Grenzwert der Polarisation bei sehr großer Zähigkeit η , v das Molekularvolumen. Nach dieser Formel berechnete F. Perrin, wie erwähnt, für Fluorescein $\tau = 4,5 \cdot 10^{-9}$ sec im Einklang mit den direkten Messungen von Gaviola.

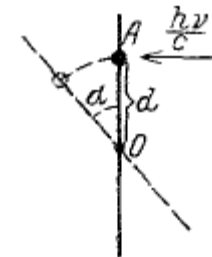


Fig. 6.

* F. Perrin, C. R. 180, 581, 1925; 182, 219.

1926; Journ. de phys. 7, 390, 1926.

**But if the fluorescence is polarized
then the measured “lifetime” will
depend on the rotation of the
molecules. Right?**

Yes

Can we get the “right” lifetime?

Yes



Alexander Jablonski
(1898-1980)

Eine Theorie der zeitlichen Abklingung des Leuchtens bei polarisierter Fluoreszenz von Farbstofflösungen.

Von A. Jabłoński in Warschau.

Mit 2 Abbildungen. (Eingegangen am 3. April 1935.)

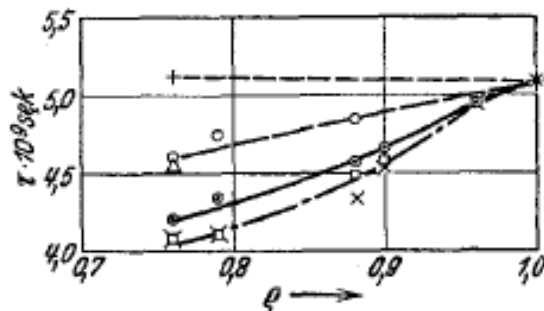
A. Jablonski, *Z. Phys.* **95**, 53 (1935).

Wir stellen also fest: *Will man direkt die Verweilzeiten der fluoreszierenden Moleküle (in Lösungen) bestimmen, so soll man (falls die Erregung mit polarisiertem Licht erfolgt) die Abklingung der unter dem Winkel $= 54,72^\circ$ gegen die Schwingungsrichtung des erregenden Lichtes schwingenden Komponente untersuchen.*

Enter Magic Angle

$$\eta = \int_0^\pi P(\varrho) \cdot f(\sigma - \varrho) d\varrho.$$

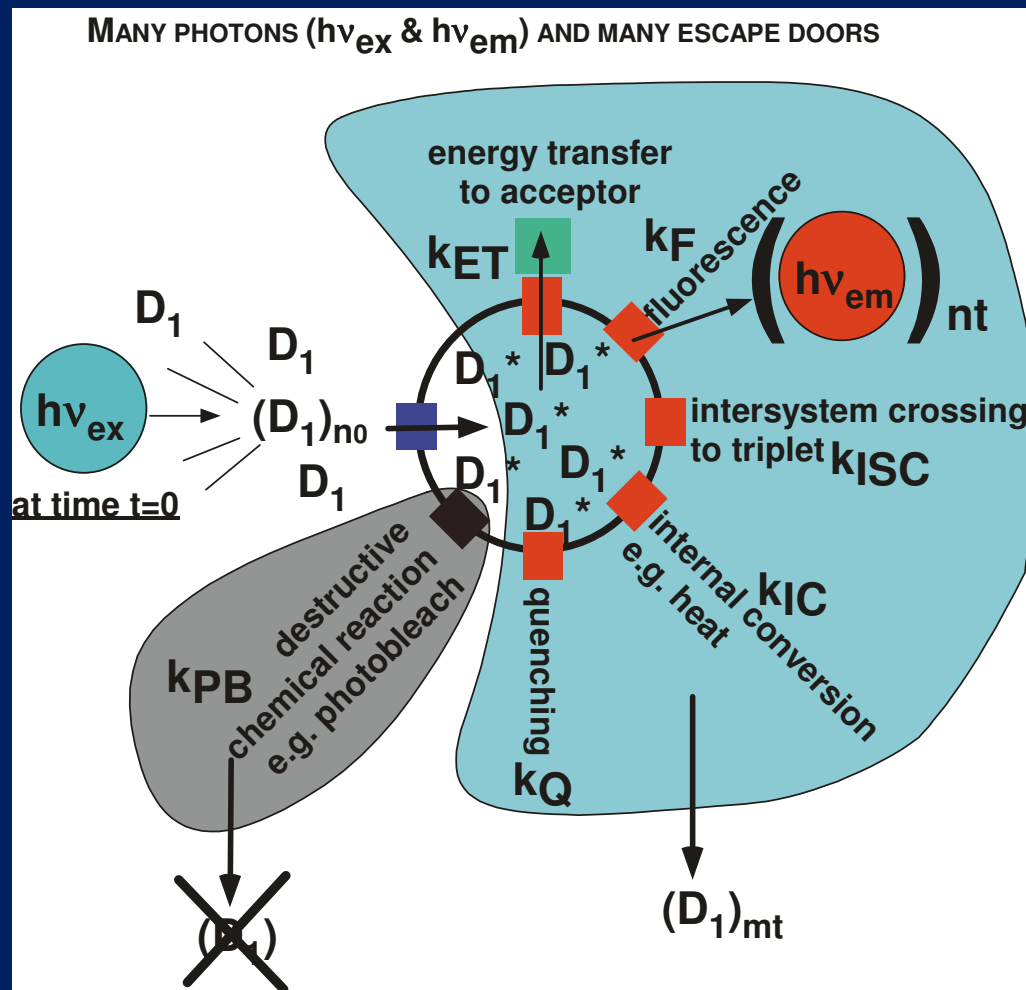
$$\left. \begin{aligned} I_{||} &= \frac{1}{3} [A_{||} + 2A_{\perp} + 2(A_{||} - A_{\perp}) e^{-\varphi t}] e^{-\frac{t}{\tau}}, \\ I_{\perp} &= \frac{1}{3} [A_{||} + 2A_{\perp} - (A_{||} - A_{\perp}) e^{-\varphi t}] e^{-\frac{t}{\tau}}. \end{aligned} \right\}$$



- + - - - - + beobachtete Werte für \bar{t}_{\perp} ,
- × - - - - × beobachtete Werte für $\bar{t}_{||}$,
- - - - - ○ berechnete Werte von $\bar{t}_{||}$, $\varrho_0 = 1/3$ vorausgesetzt,
- - - - - □ berechnete Werte von $\bar{t}_{||}$, $\varrho_0 = 0,1$ vorausgesetzt,
- △ beobachteter Wert für die rotationsunbeeinflussbare Komponente,
- - - - - ○ berechnete Werte für dieselbe Komponente.

Fig. 2. Abhängigkeit der Abklingzeiten der Fluoreszenz von Fluorescein in Glycerin + Wasser-Lösungen von verschiedenem Wassergehalt von dem Depolarisationsgrade ϱ der Fluoreszenz. $\varrho = 1$ entspricht der reinen wässrigen Lösung.

So, the lifetime depends on many things



lifetime of the excited state:

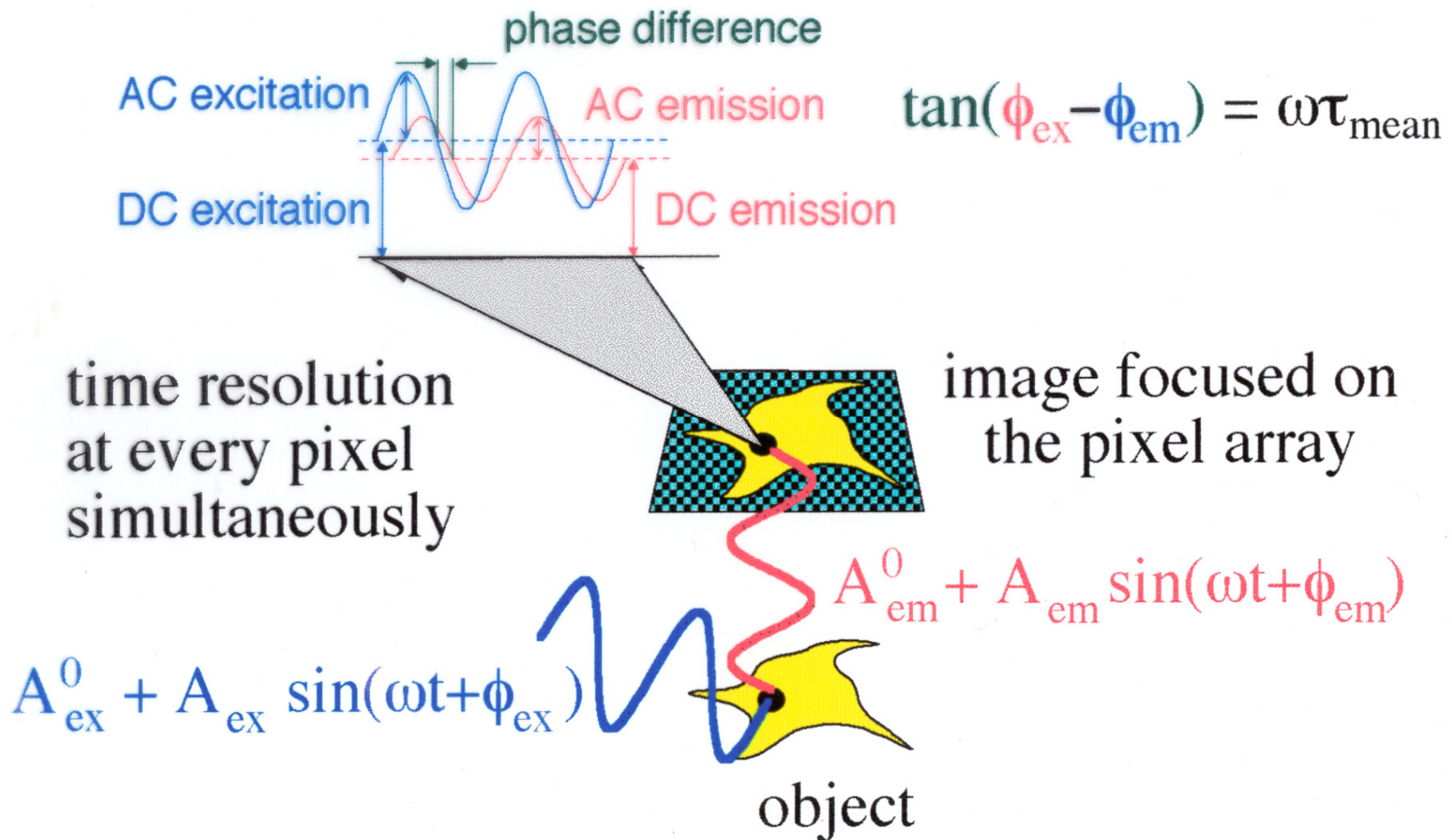
$$1/\tau = k_T + k_F + k_{ISC} + k_{IC} + k_Q + k_{PB} = \sum_j k_j$$

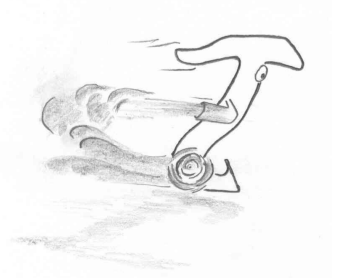
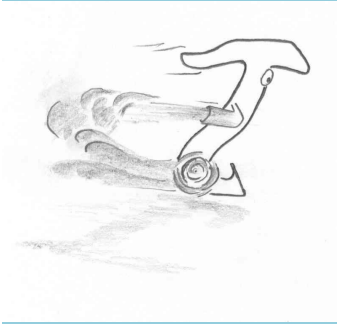
$$\text{Quantum yield of the } i^{\text{th}} \text{ process} = \frac{k_i}{\sum_j k_j}$$

SO, now we seem all set.

BUT.....

We want to measure fluorescence lifetimes in a fluorescence image at every location of the cell.



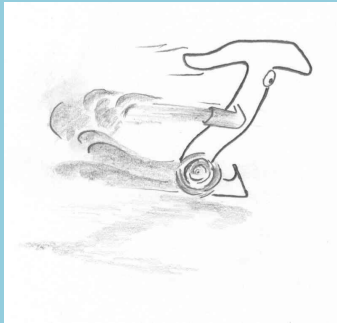


I see three exponentials!



Uh-Oh! I see a whole herd of 'em!

exponentials!



What now?



Lifetimes in images are not so simple!
We need some help!

Let's start with

Model Independent Analysis

Some different ways to
parameterize lifetime-resolved
data

$$1/(1 + j\omega t) = M_i \left[\cos(\phi_{i,\omega}) + j \sin(\phi_{i,\omega}) \right]$$

$$x = M_i \cos(\phi_{i,\omega}) \text{ and } y = M_i \sin(\phi_{i,\omega})$$

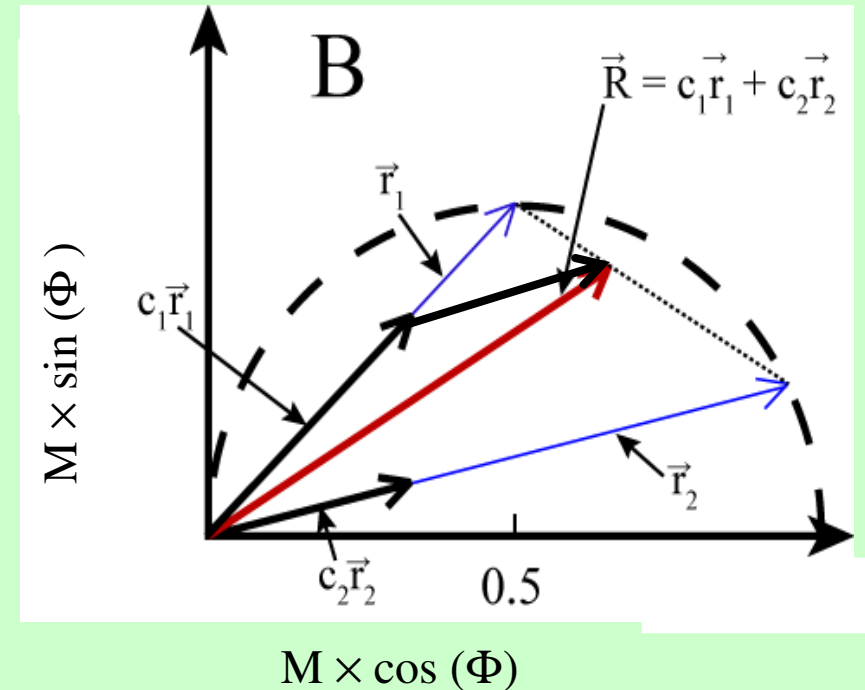
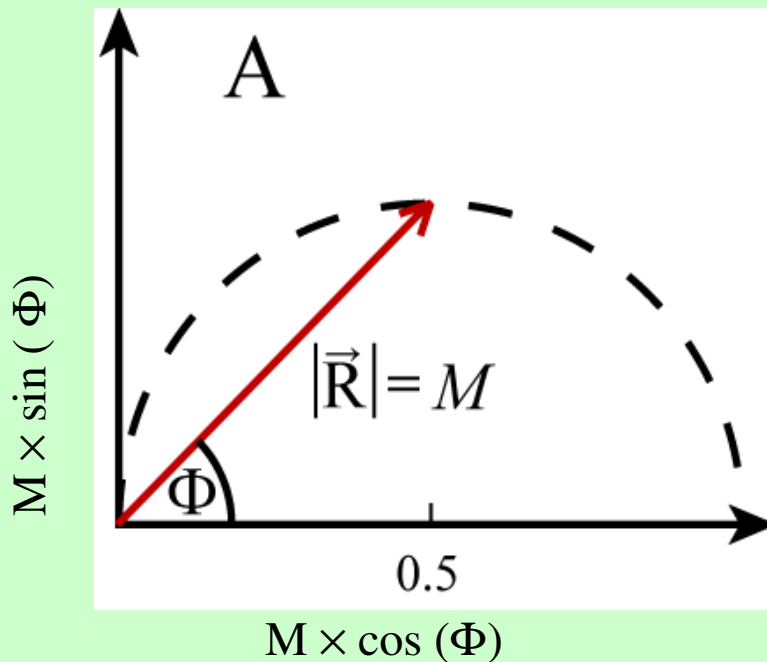
$$j = \sqrt{-1}$$

Frequency domain lifetime measurement

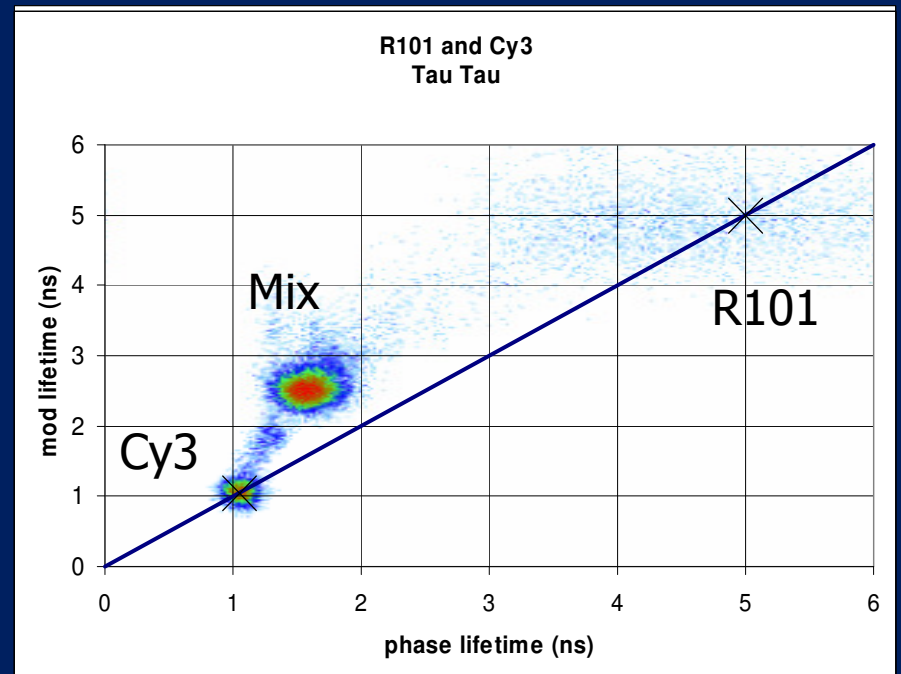
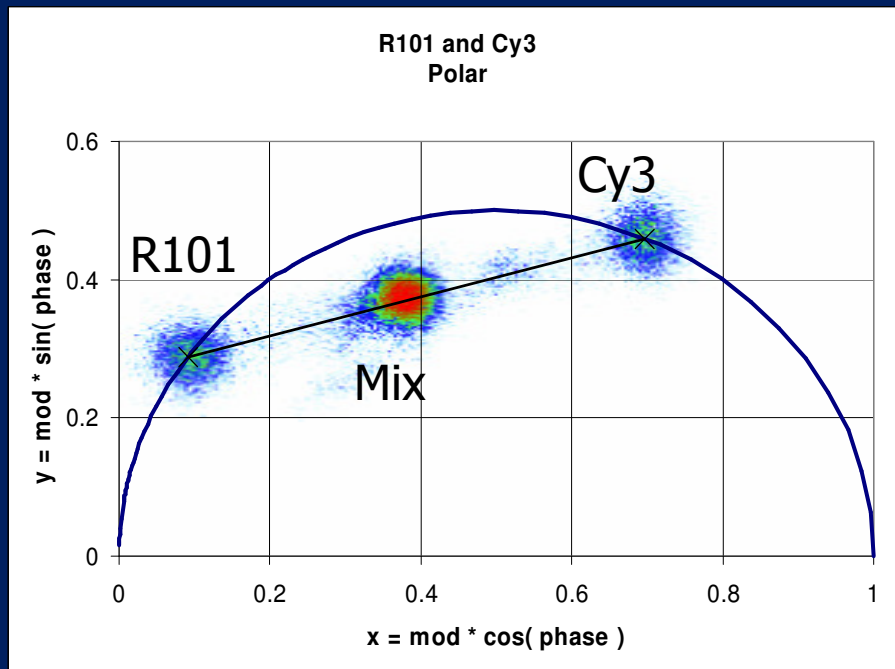
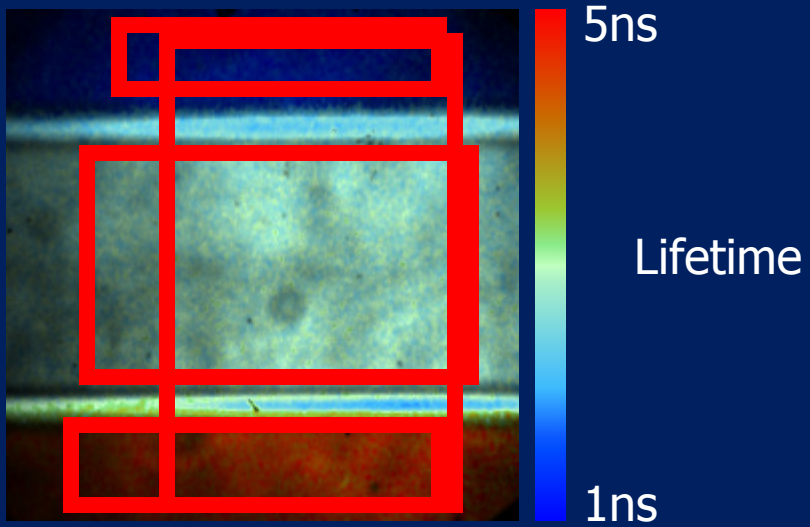
Data analysis with a **polar plot** representation

Demodulation = $M = \frac{b/B}{a/A} = \frac{1}{\sqrt{1 + (\omega\tau_M)^2}}$

Phase shift = $\Phi = \tan^{-1}(\omega\tau_\Phi)$



good for any signal $\propto \frac{1}{1 + i\omega\tau}$ (for instance dielectric dispersion)

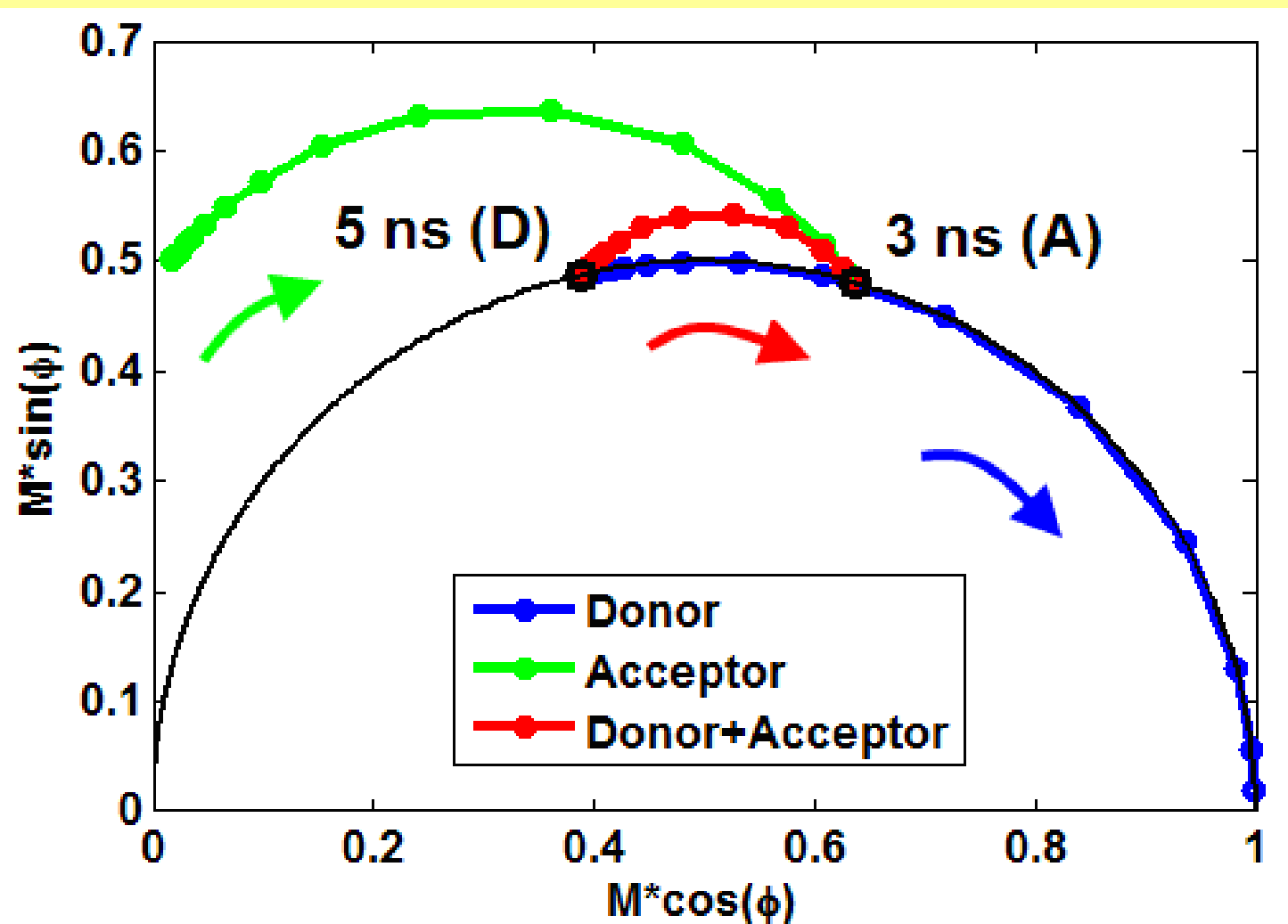


Observing the fluorescence of:

Product species of an excited state reaction

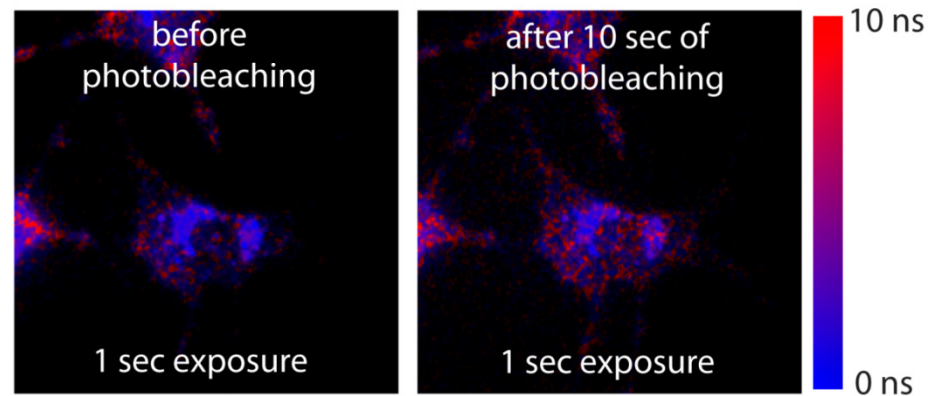
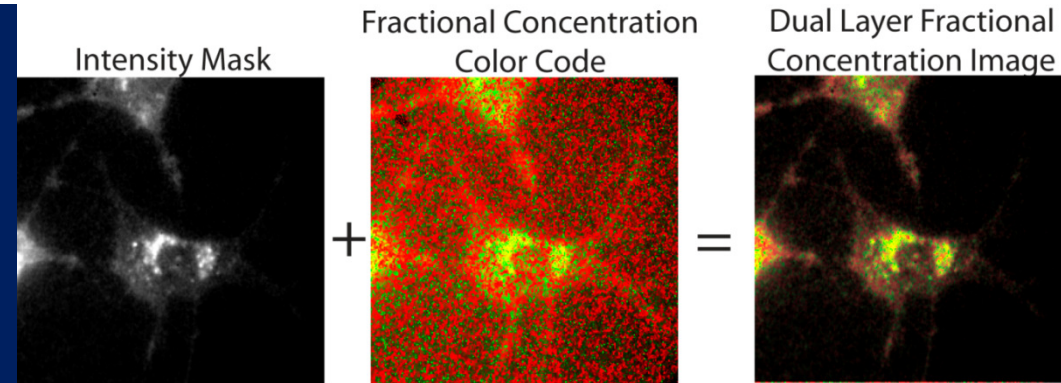
Product and directly excited species

Directly excited species



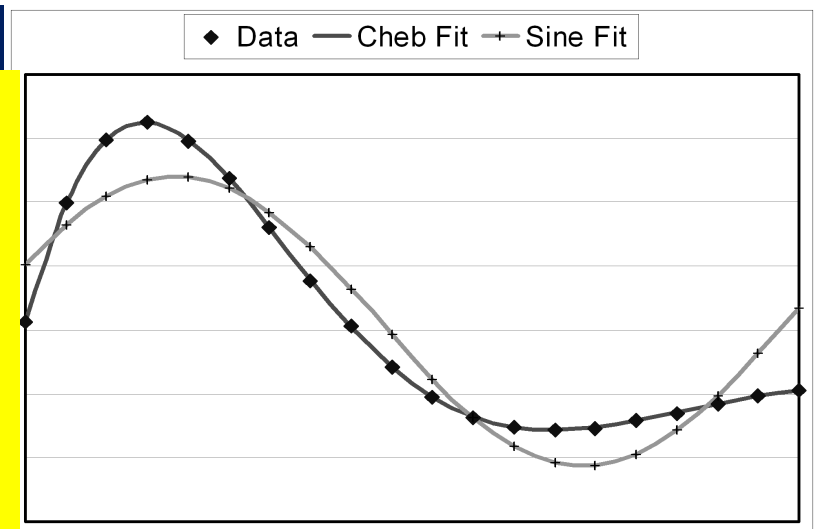
Of course we have to be sure that we get rid of artifacts

Chebyshev fits
- very fast
single pass
not interactive

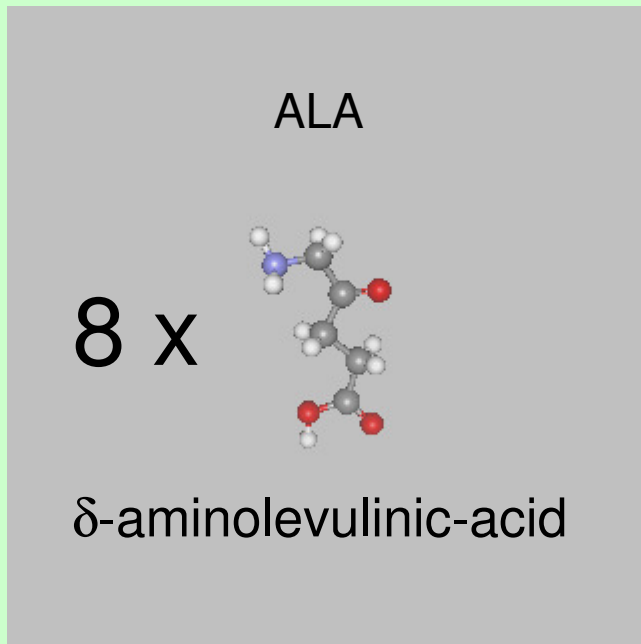


Correction for photobleaching

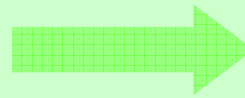
Malachowski, G.C., Clegg, R.M., and Redford, G.I., "Analytic solutions to modelling exponential and harmonic functions using Chebyshev polynomials: fitting frequency-domain lifetime images with photobleaching," *Journal of Microscopy* 228(3), 282-295 (2007)



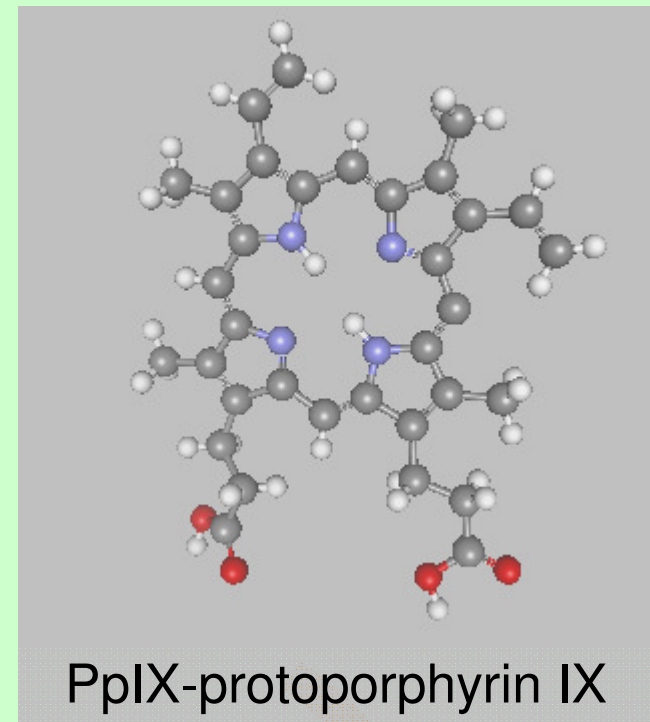
Photodiagnosics and phototherapy



Normal cellular

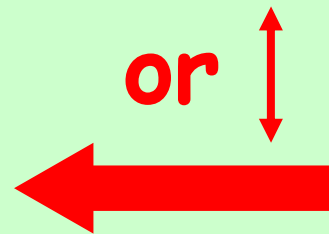


heme synthesis

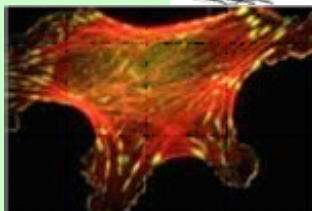
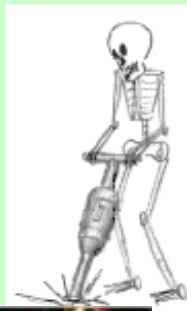
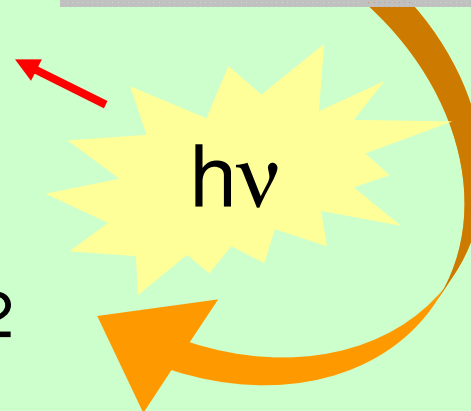


Fluorescence

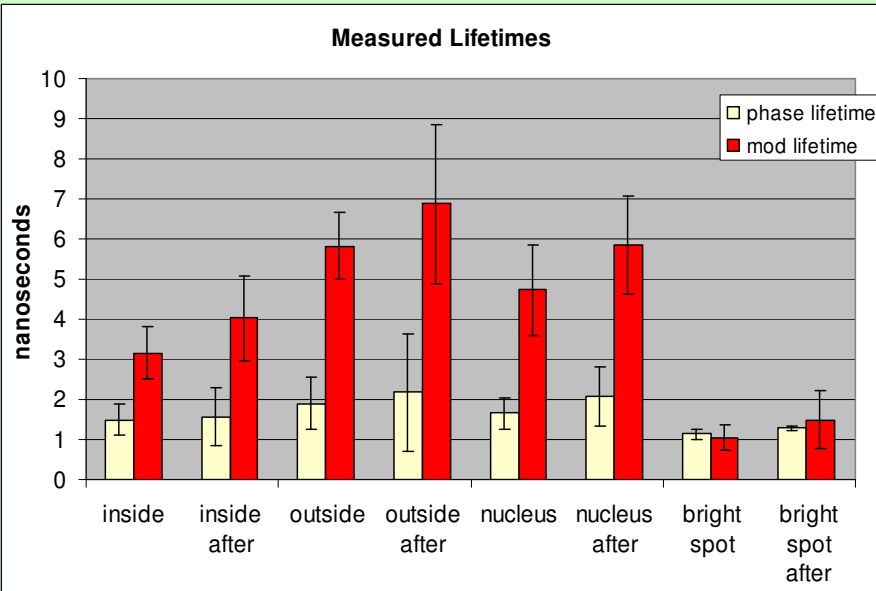
or



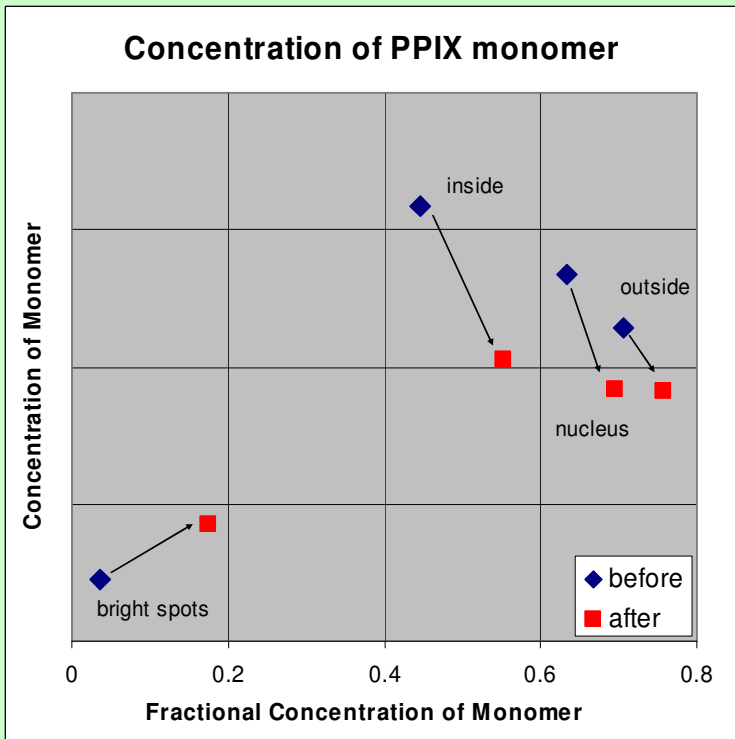
$^1\text{O}_2$



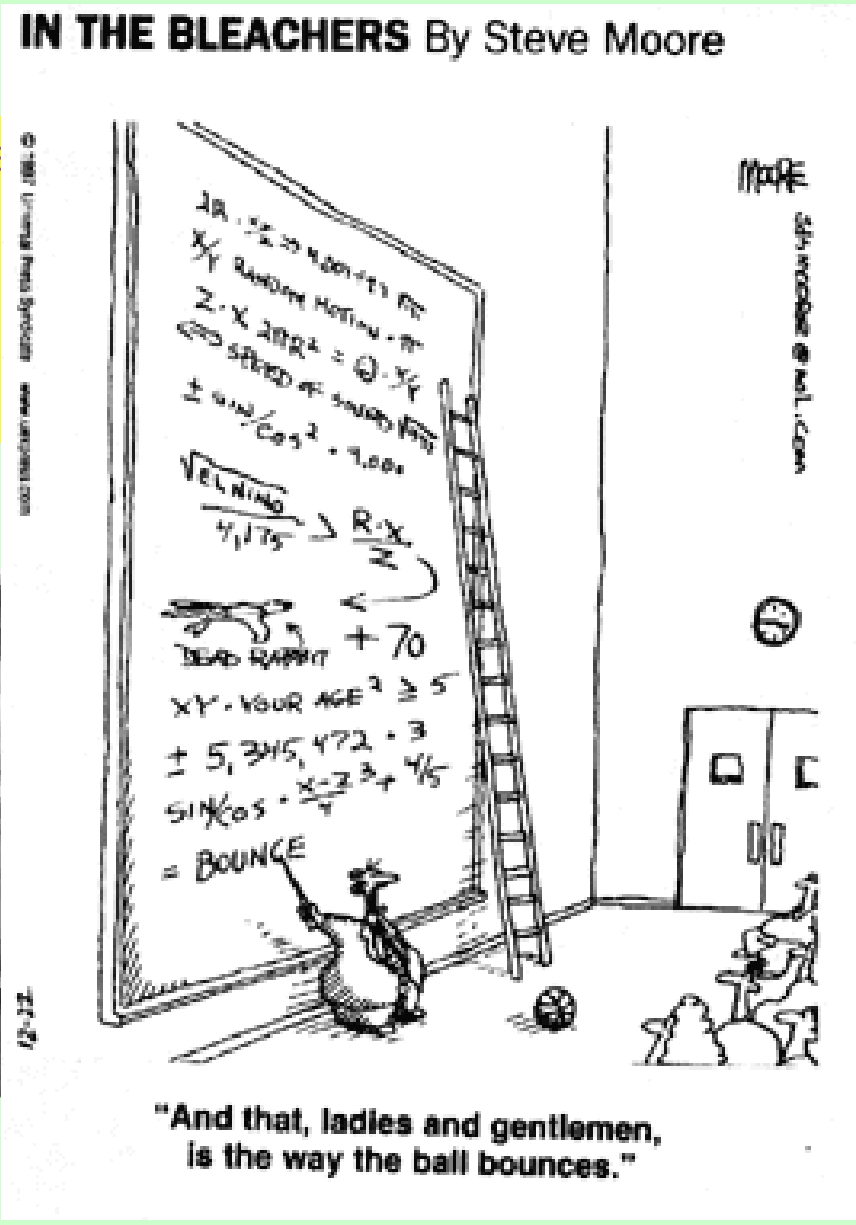
The monomer of PpIX forms ROS and is used for phototherapy



We

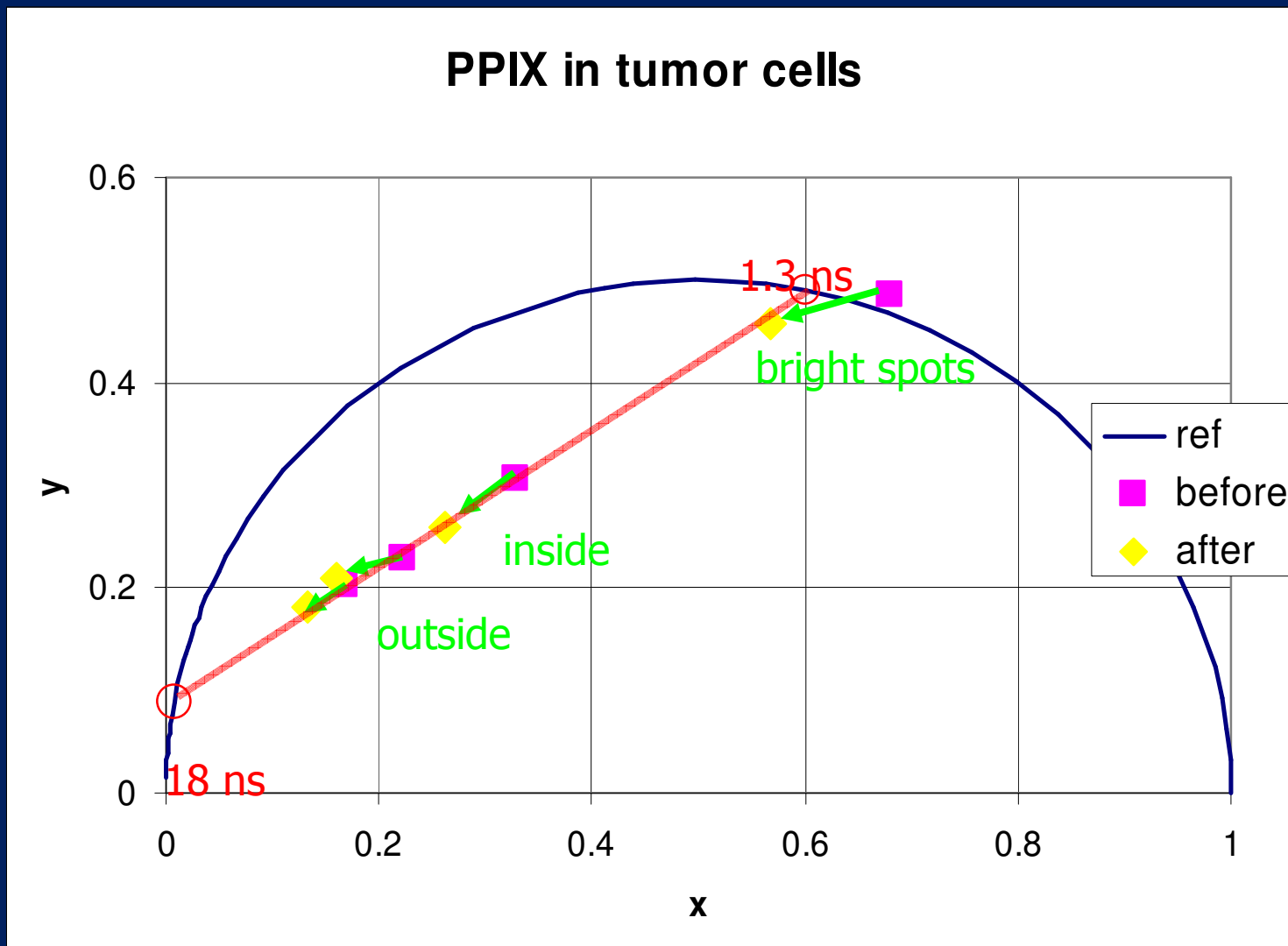


But sometimes that complicates things



Polar Plot analysis

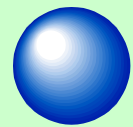
Fluorescence intensity and dynamic response change upon illumination



Protoporphyrin IX

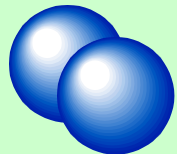
Lifetime

Location



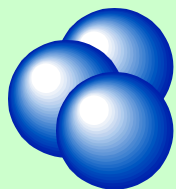
Monomer

17-18 ns



Dimer

~2 ns

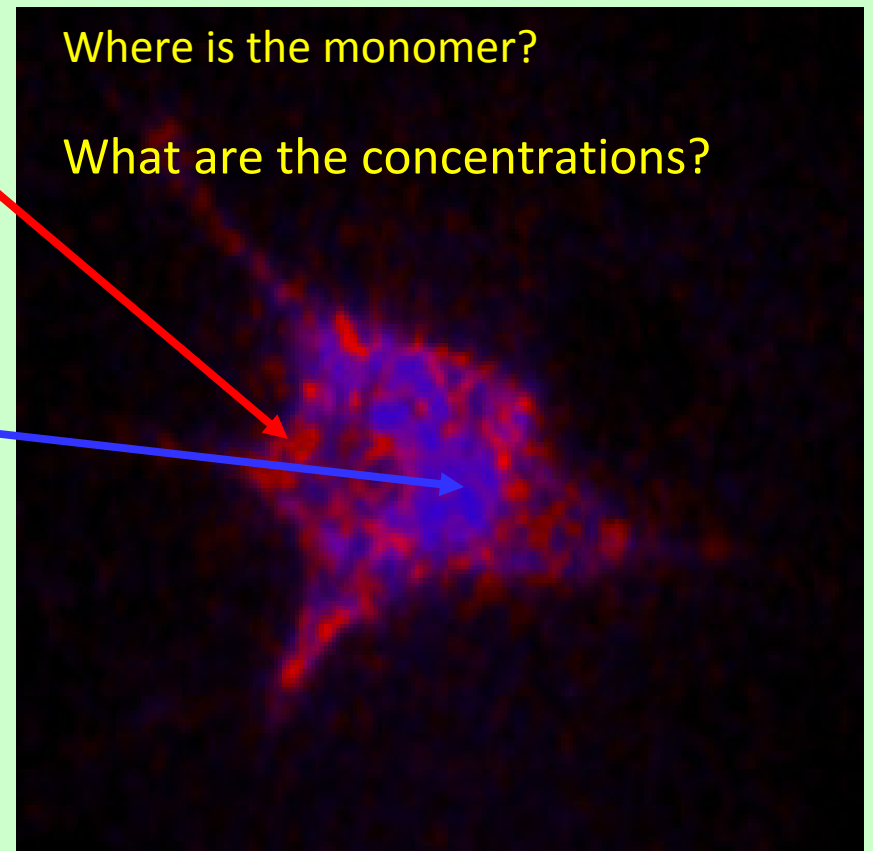


Higher
Multimers

<2 ns

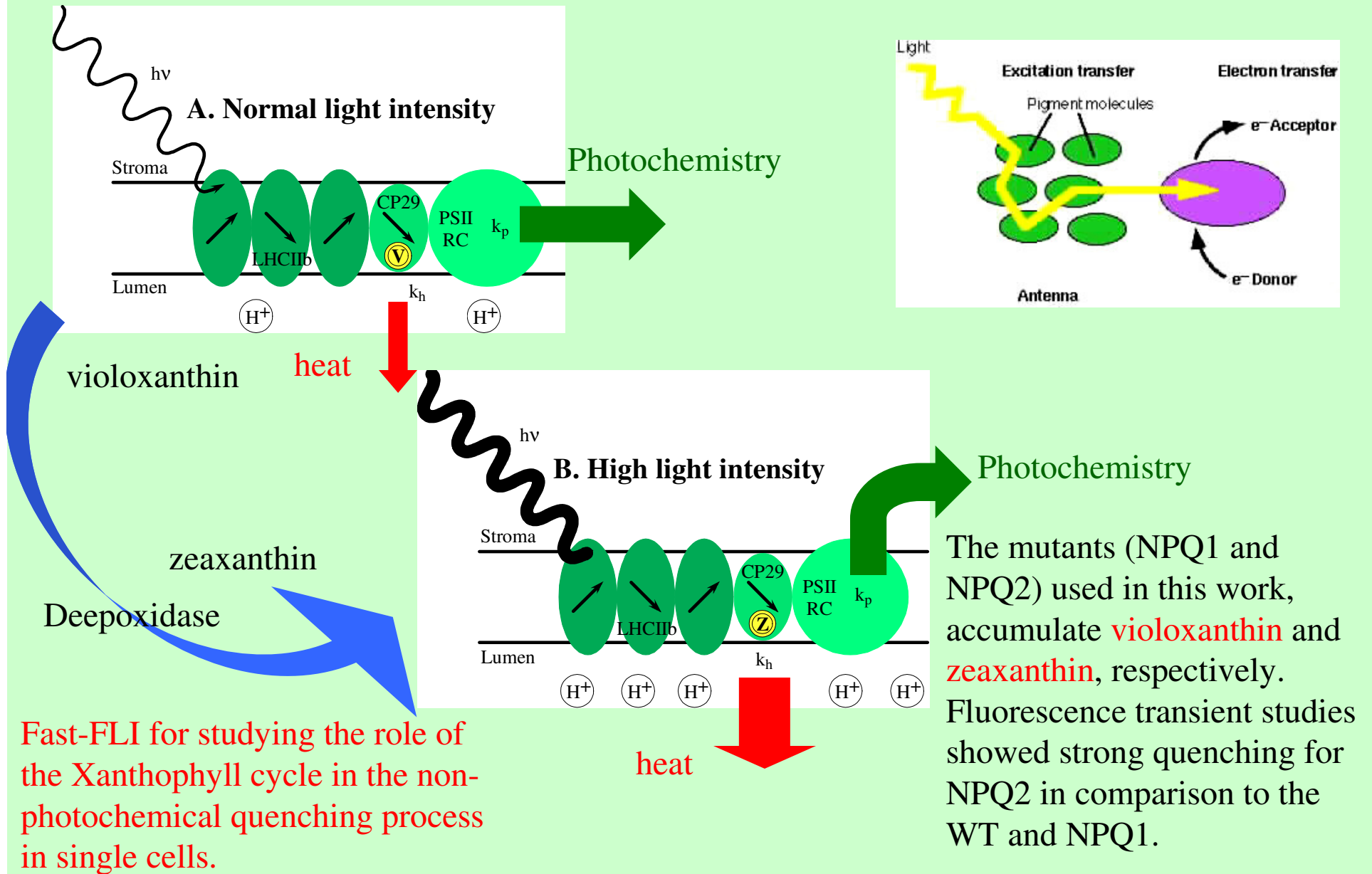
Where is the monomer?

What are the concentrations?



Monomers and Multimers \approx same spectra

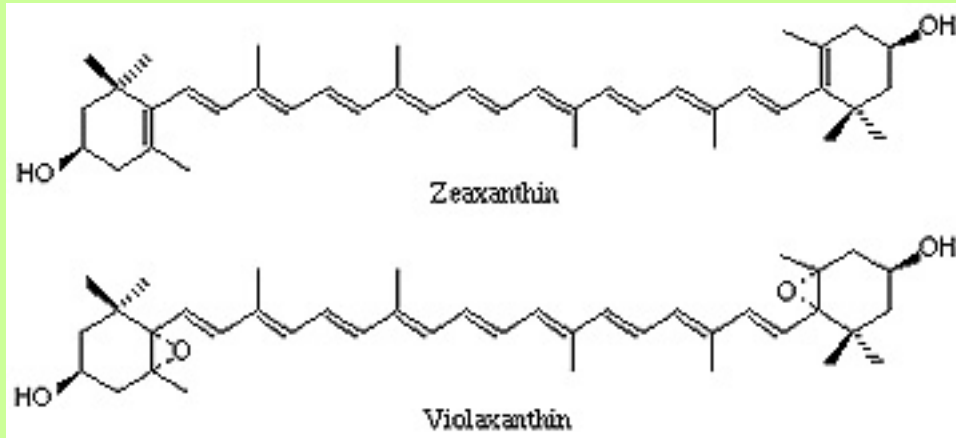
Real-Time Fluorescence Lifetime-Resolved Images of individual cells of Wild Type and NPQ mutants of *Chlamydomonas reinhardtii*



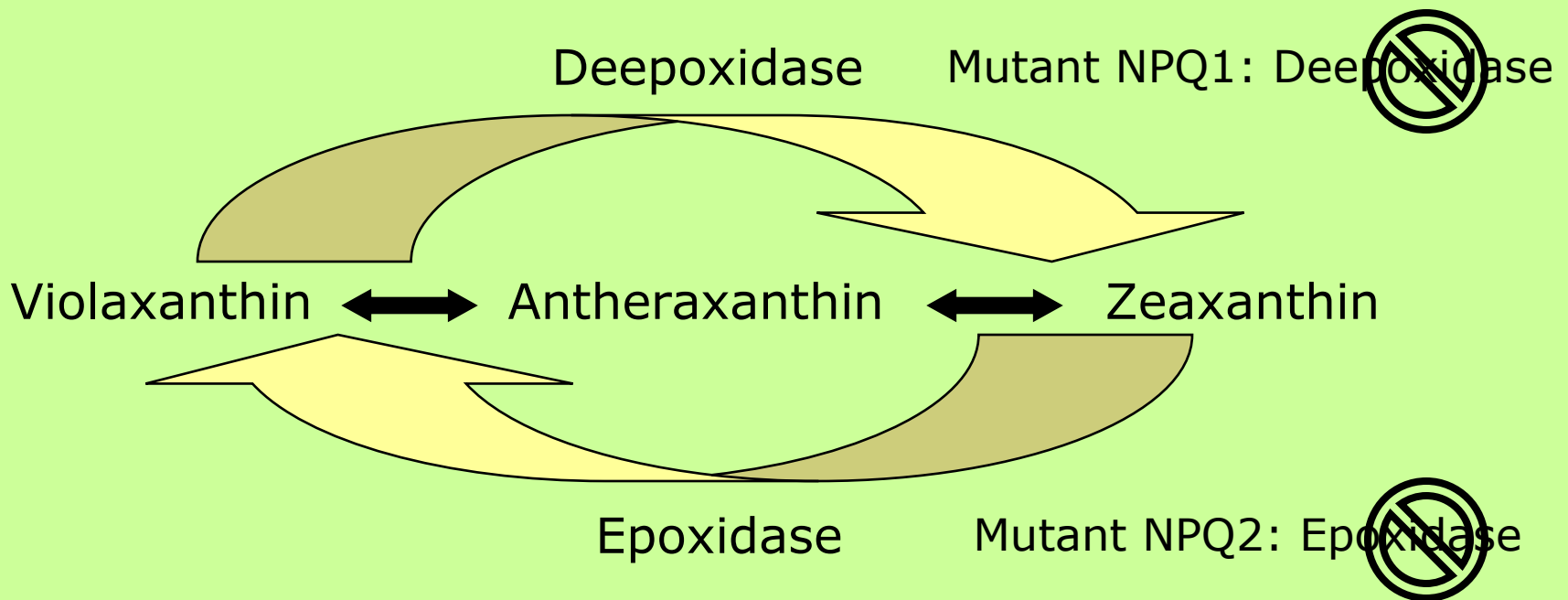
Fast-FLI for studying the role of the Xanthophyll cycle in the non-photochemical quenching process in single cells.

NPQ mutants of the green alga *Chlamydomonas reinhardtii*

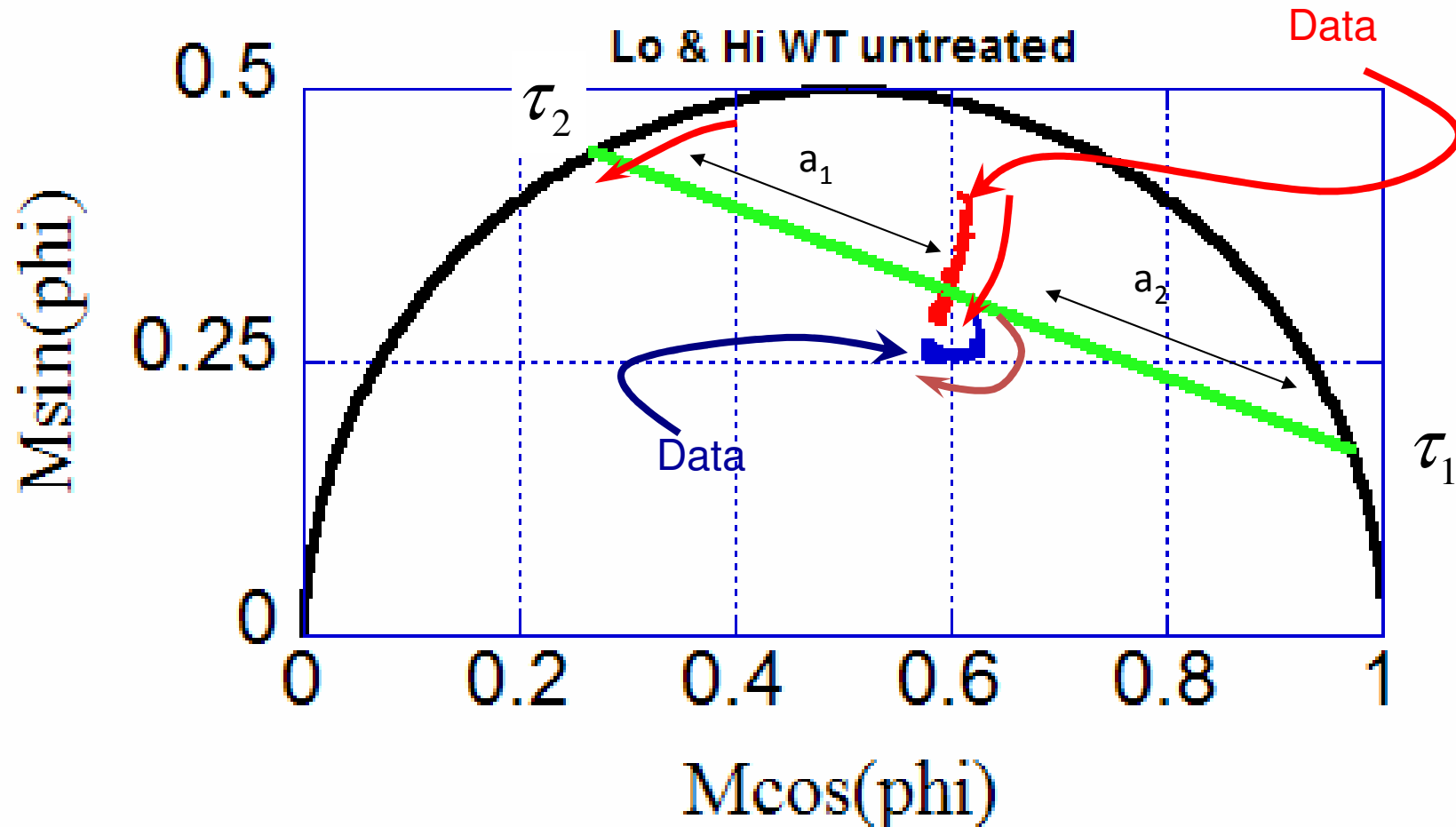
Mutants from Krishna K. Niyogi



The Xanthophyll cycle



Lifetime lever - long time biochemical changes in the cells

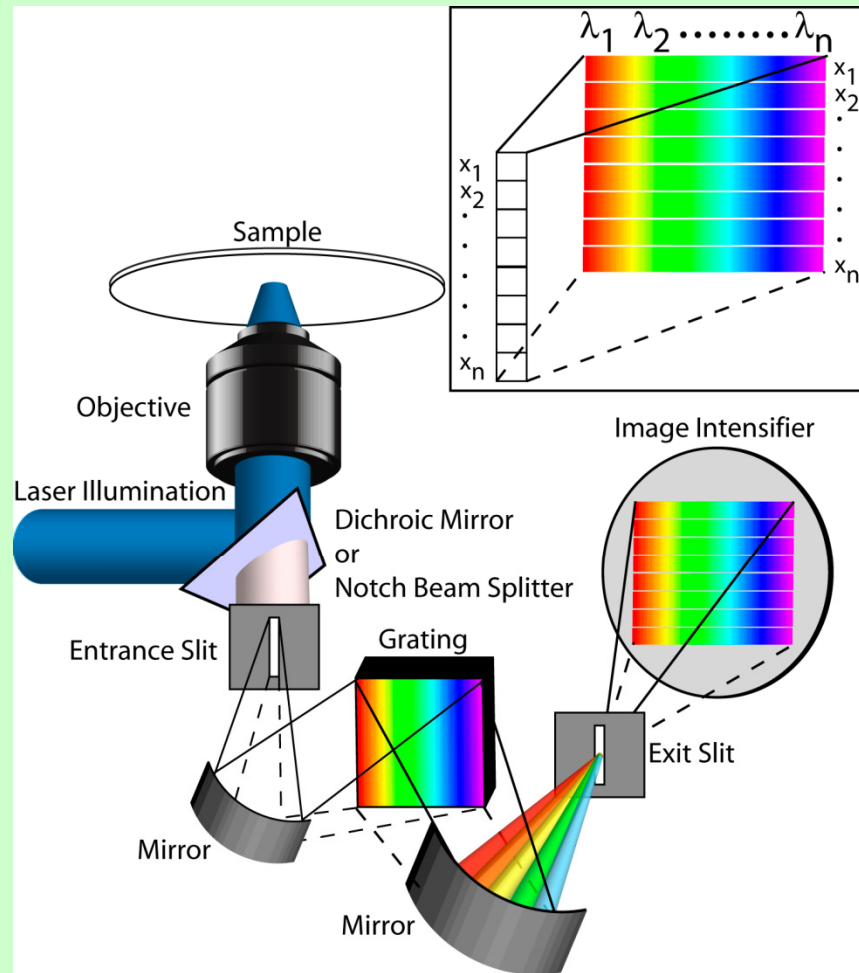


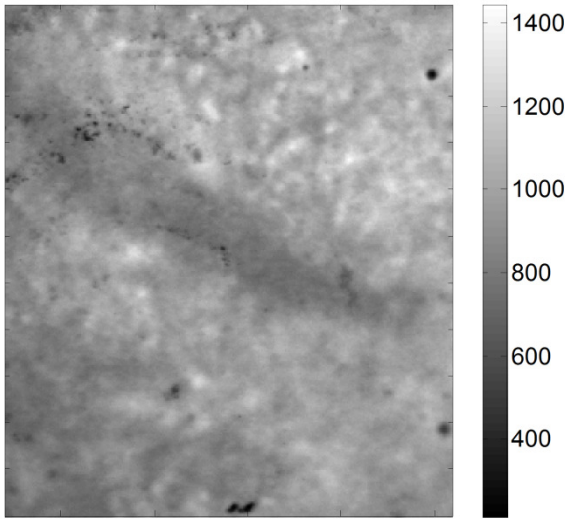
$$\text{Signal} \neq a_1 e^{-t/\tau_1} + a_2 e^{-t/\tau_2}$$

Two lifetime pools of Chlorophyll Molecules

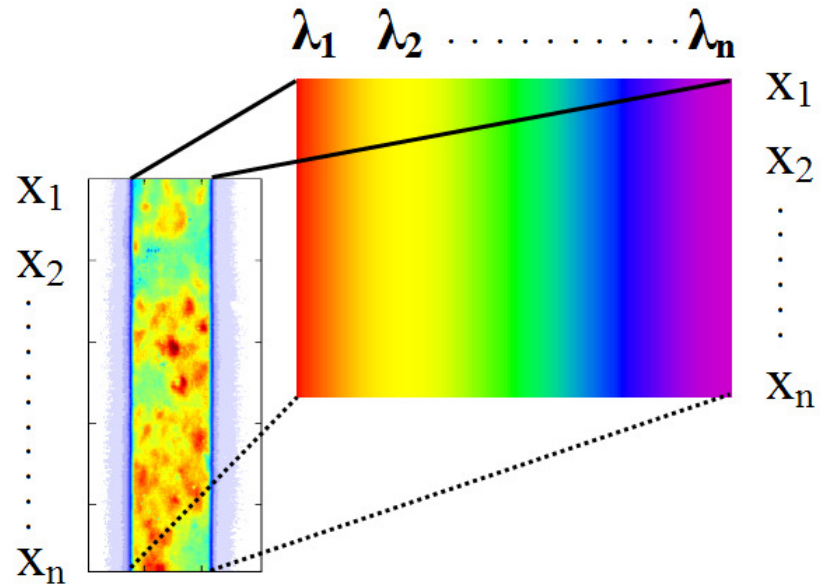
Holub, O., Seufferheld, M.J., Gohlke, C., Govindjee, Heiss, G.J., and Clegg, R.M.,
"Fluorescence lifetime imaging microscopy of Chlamydomonas reinhardtii: non-

Spectral FLIM

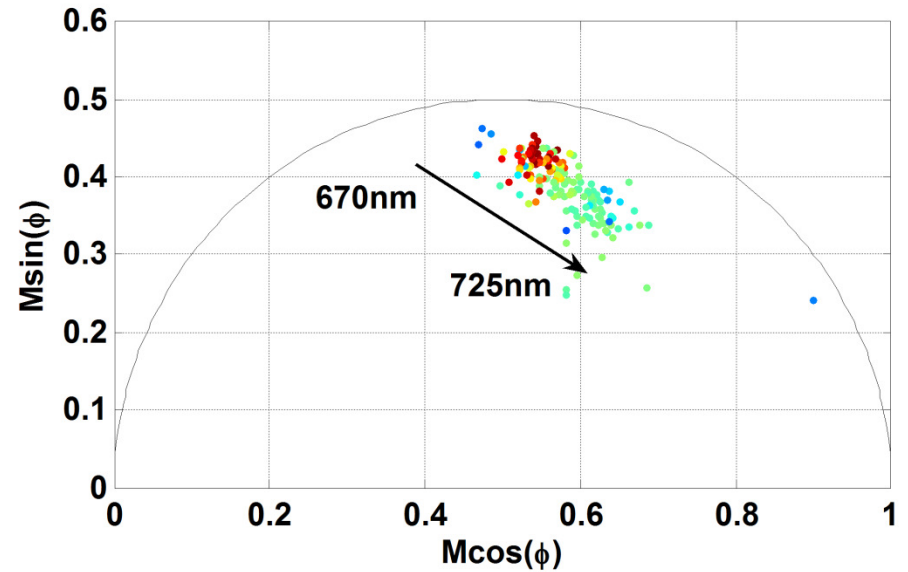
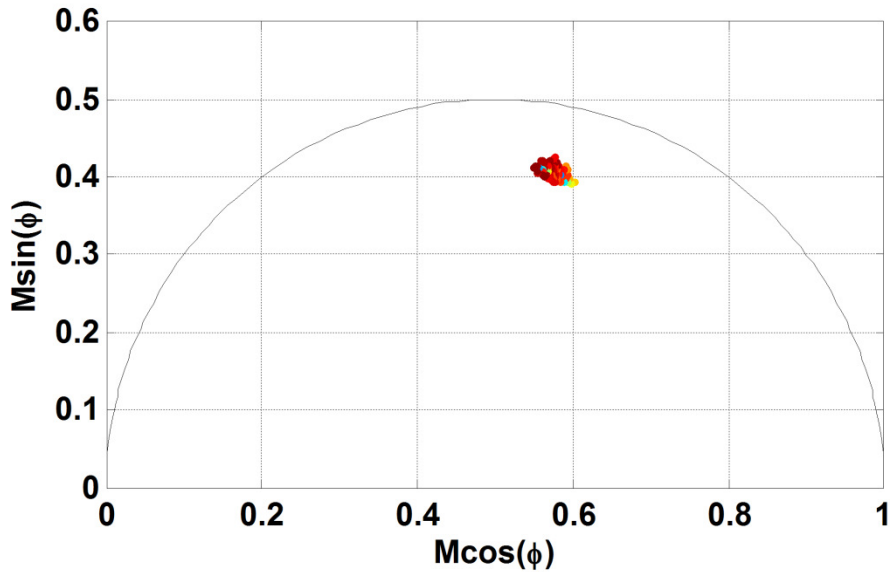




Avocado Leaves



Different wavelengths - different lifetimes



Combining morphology + lifetime resolution

Localized spatial frequencies

$$w(x, y, a) = \iint g\left(\frac{x-x'}{a}, \frac{y-y'}{a}\right) f(x', y') dx' dy'$$

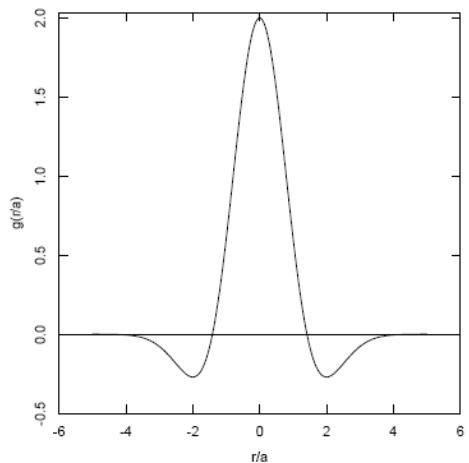


FIG. 1.—Mexican hat-generating wavelet $g(r/a)$

Wavelets + Denoising

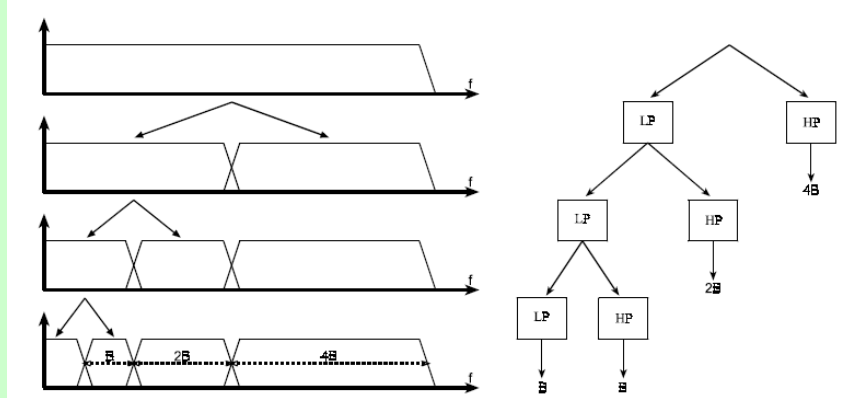
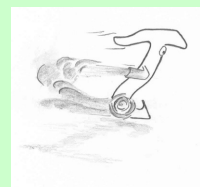
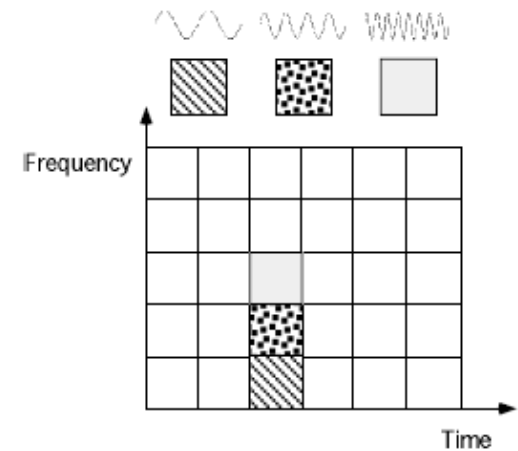


Figure 4
Splitting the signal spectrum with an iterated filter bank.

Chasing lifetimes in the morphology and noise



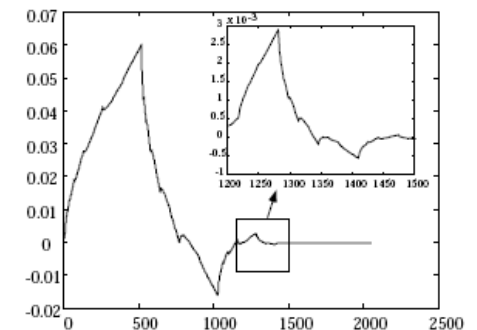
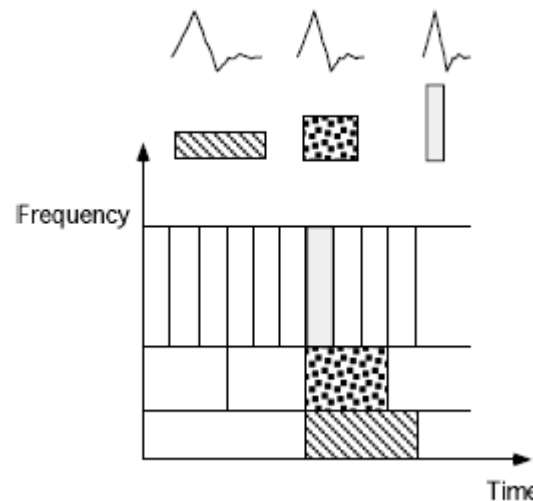
$$a_0 + \sum_{k=1}^{\infty} (a_k \cos kx + b_k \sin kx)$$



Fourier basis functions, time-frequency tiles, and coverage of the time-frequency plane.

$$\Phi_{(s,l)}(x) = 2^{-\frac{s}{2}} \Phi(2^{-s}x - l)$$

$$W(x) = \sum_{k=-1}^{N-2} (-1)^k c_{k+1} \Phi(2x + k)$$



Daubechies mother wavelet

Daubechies wavelet basis functions, time-frequency tiles, and coverage of the time-frequency

WHAT DO SOME WAVELETS LOOK LIKE?

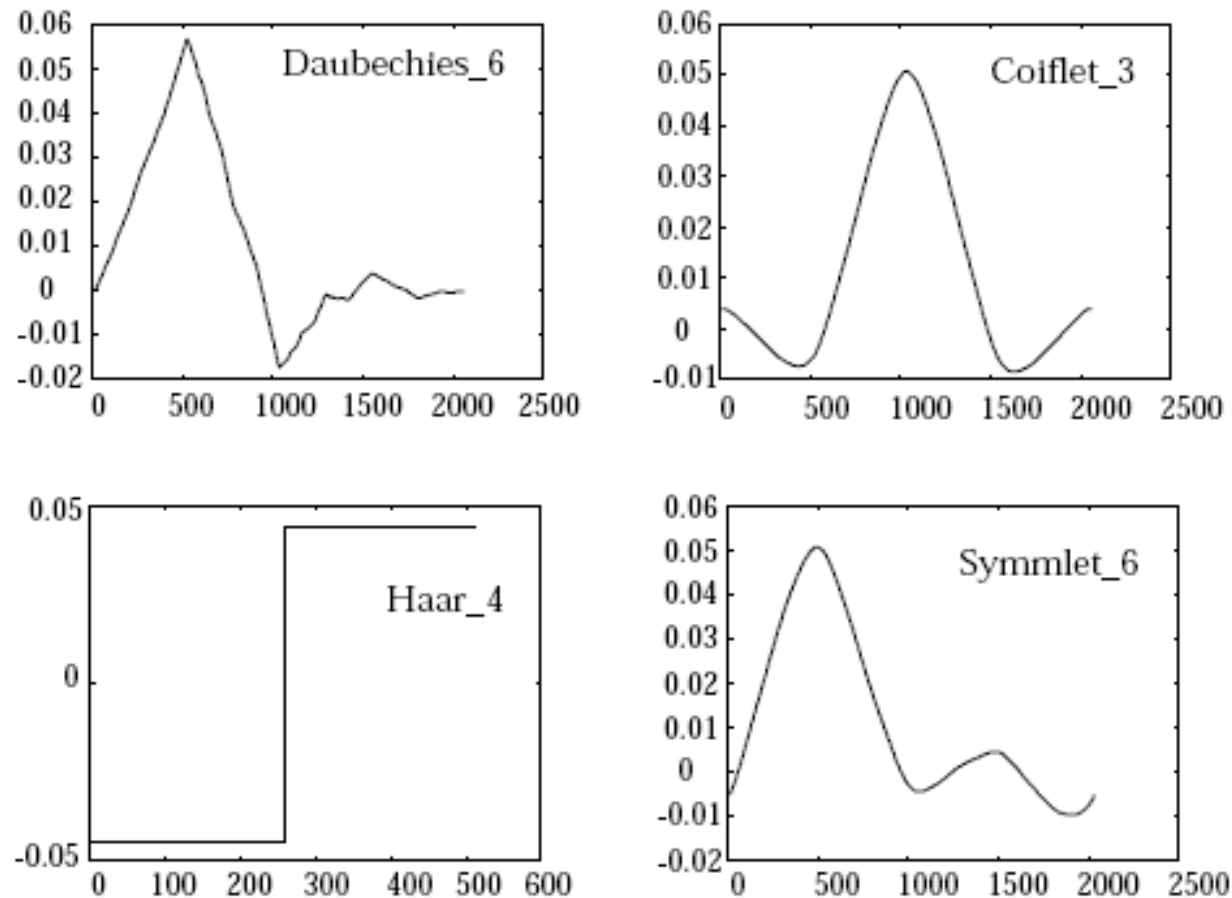
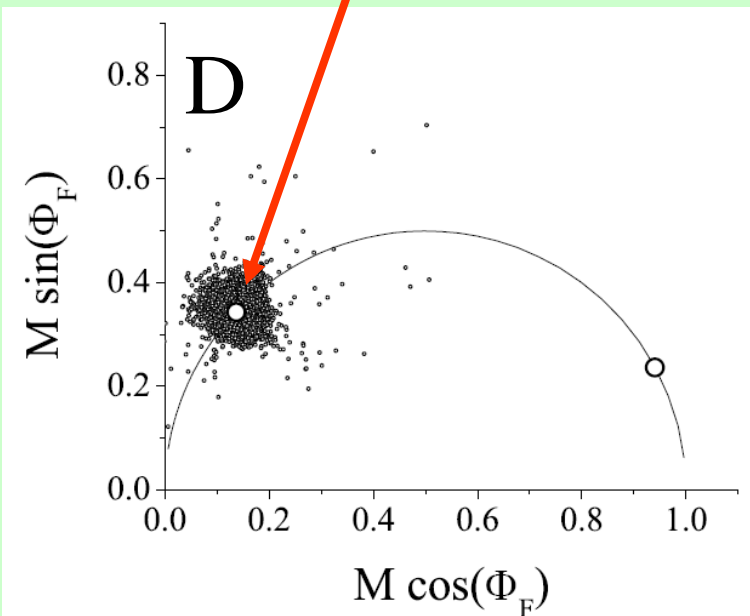
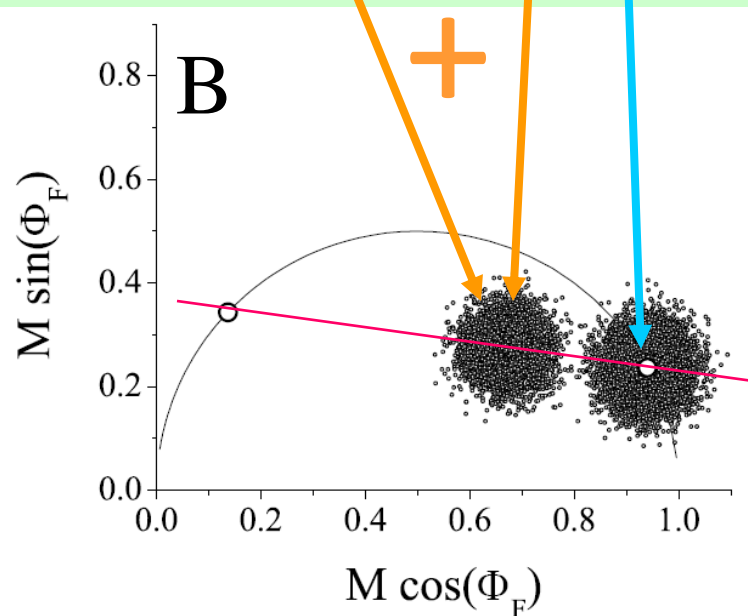
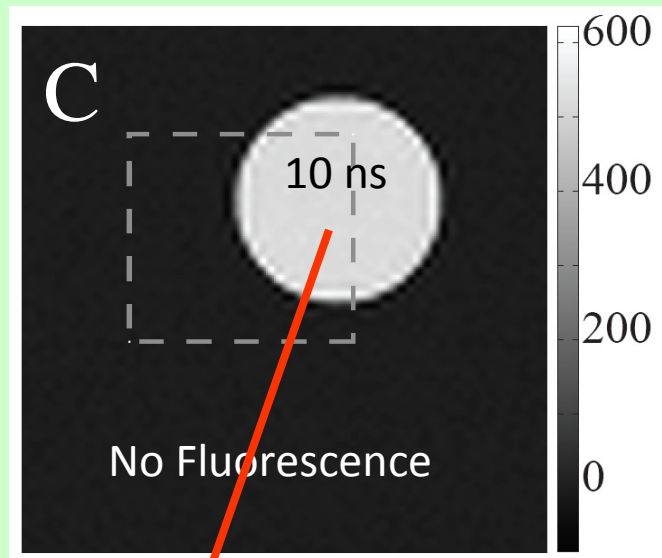
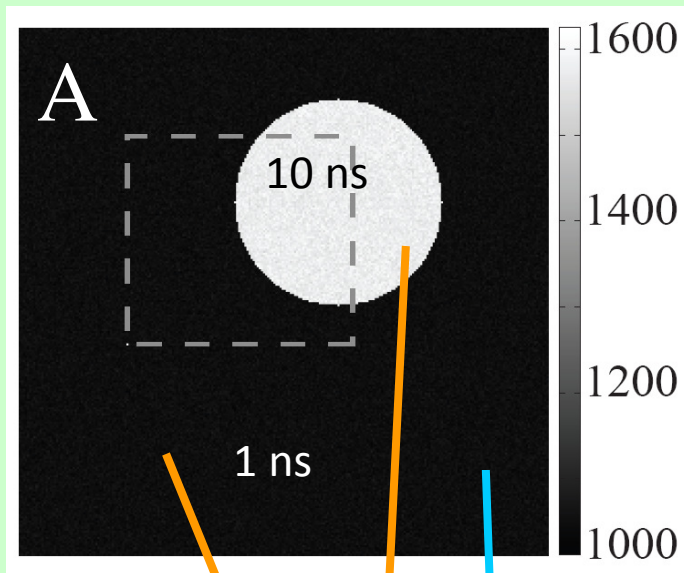


Fig. 4. Several different families of wavelets. The number next to the wavelet name represents the number of vanishing moments (A stringent mathematical definition related to the number of wavelet coefficients) for the subclass of wavelet. These figures were generated using WaveLab.

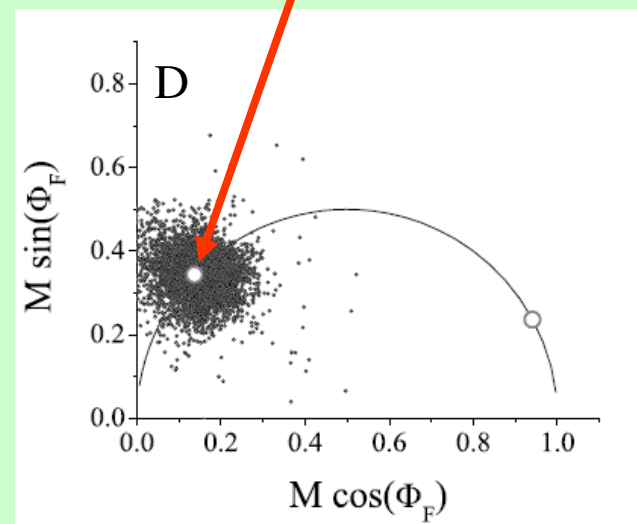
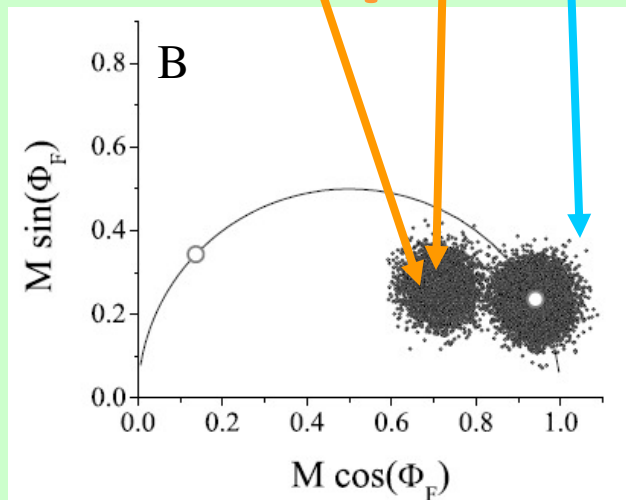
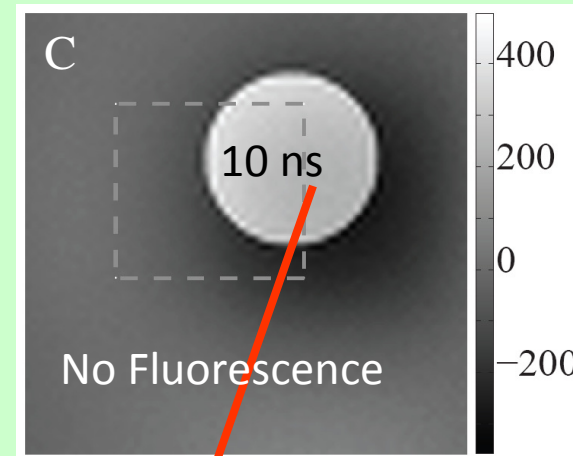
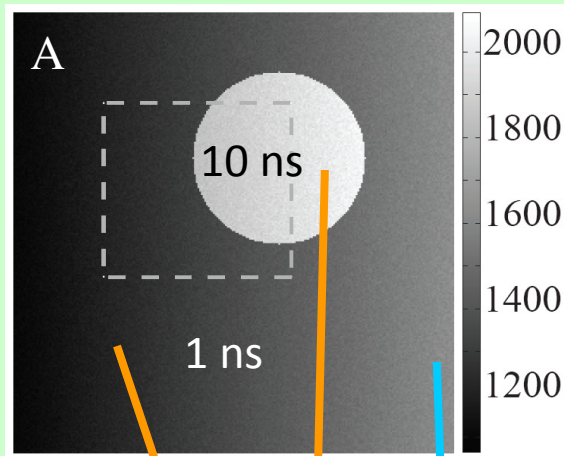
combination of wavelet analysis and FLIM with simulated data

Background constant

The wavelet analysis has completely removed the background contribution



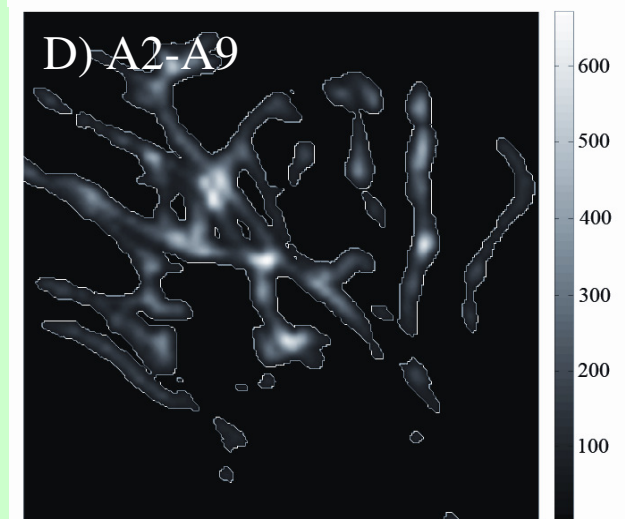
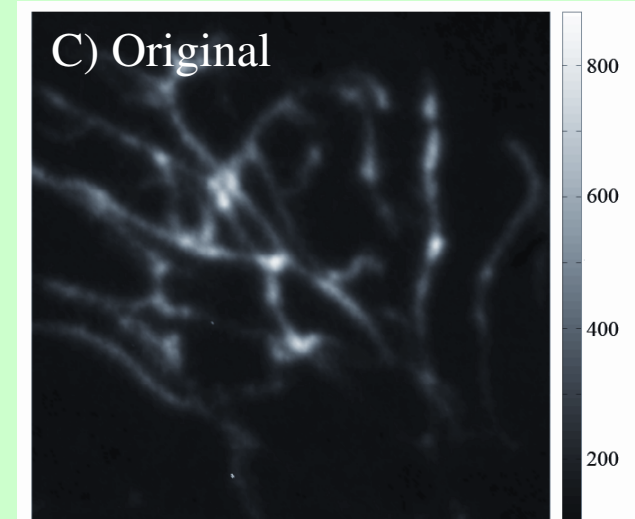
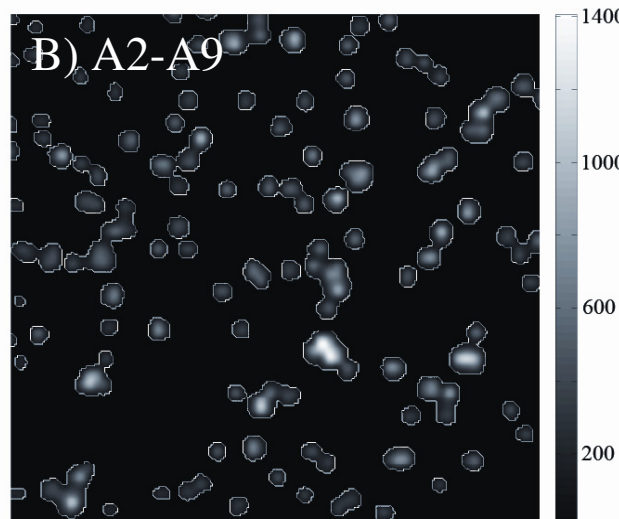
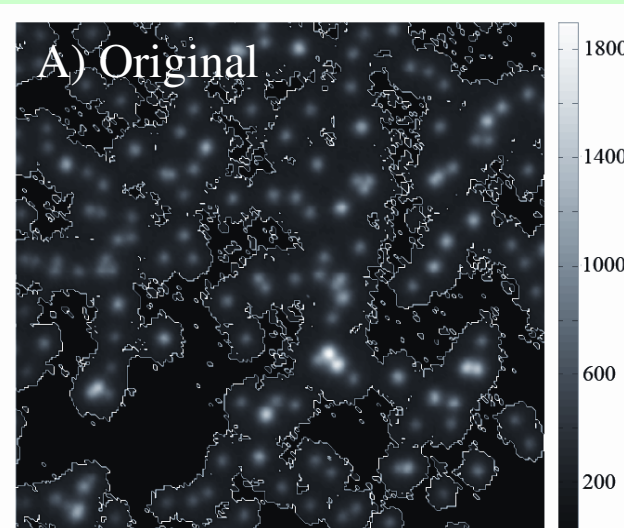
Combination of wavelet analysis and FLIM with simulated data
The background in this simulated data is increasing in amplitude with a constant gradient from left to right.



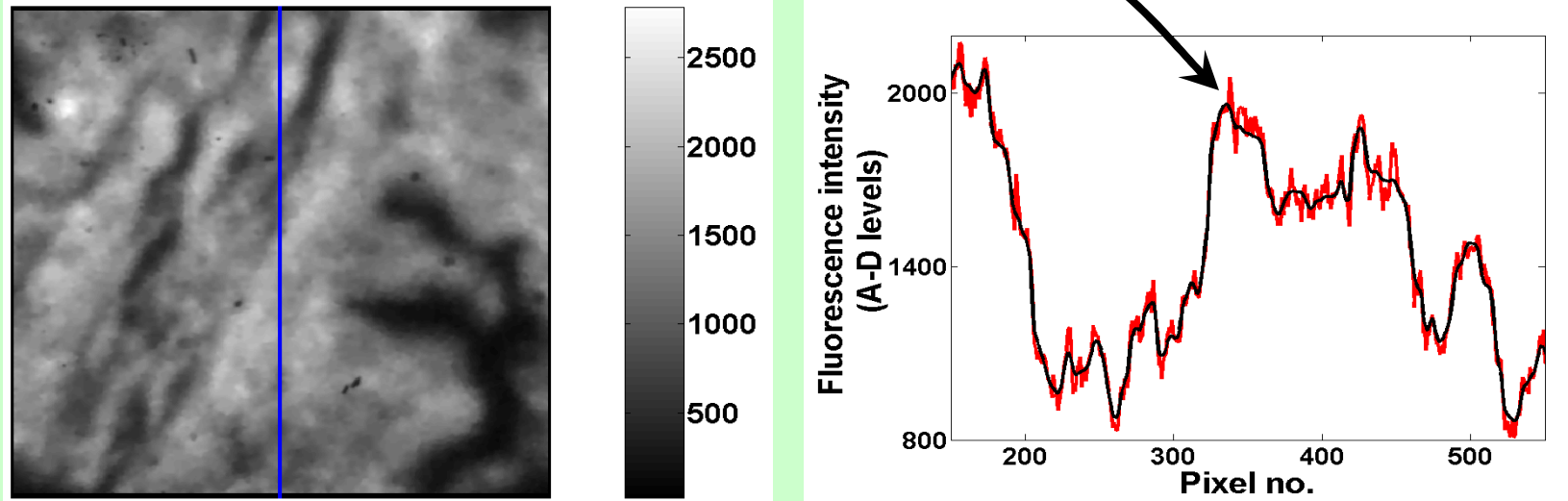
Finding morphology using wavelets

Background subtraction using wavelet on the **fluorescent beads** image (A and B) and on the **dendrites in a *Drosophila melanogaster* larva** expressing membrane-tagged GFP (C and D). The original images (A and C) and the edited image analyzed with wavelet (B and D) are compared.

final images are reconstructed by the 'wrcoef2' function from the **difference** in the approximation **data level 2** (containing both high and low spatial frequency components) and **level 9** (containing mostly low spatial frequency component)



Variance stabilized Gaussian Denoising



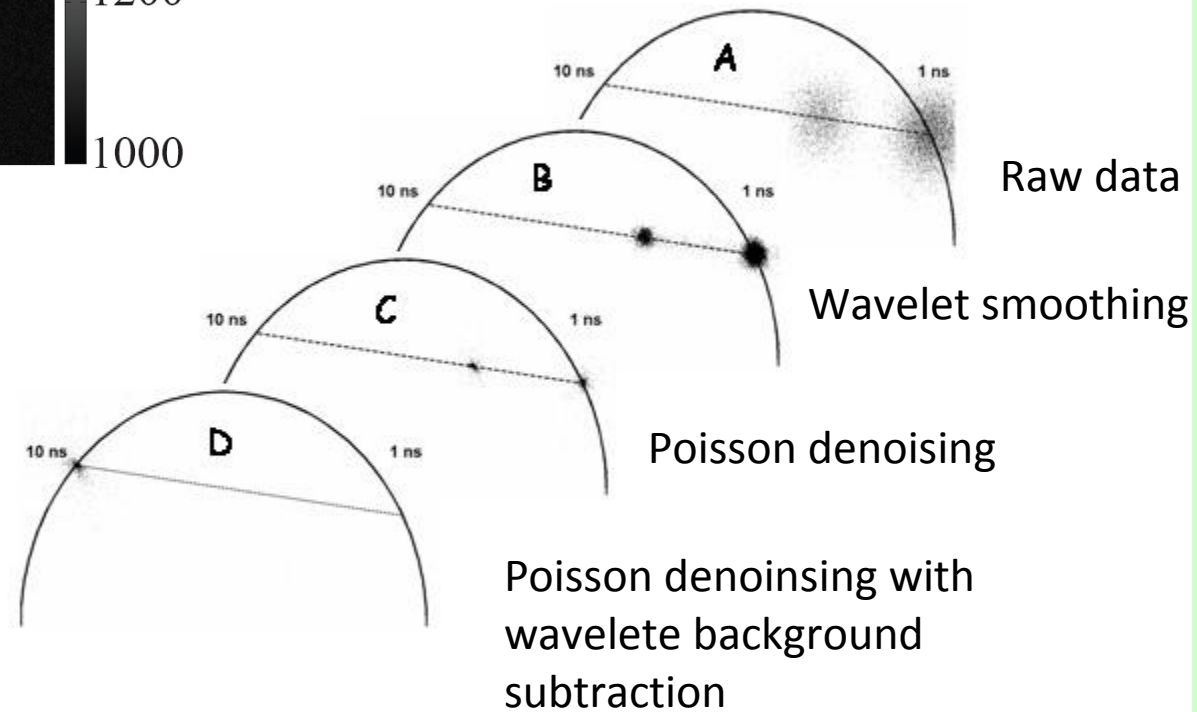
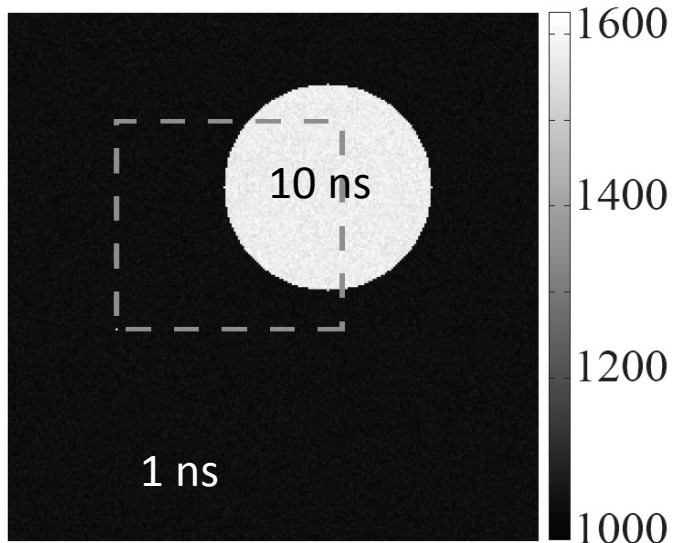
Red – raw data; Black – denoised

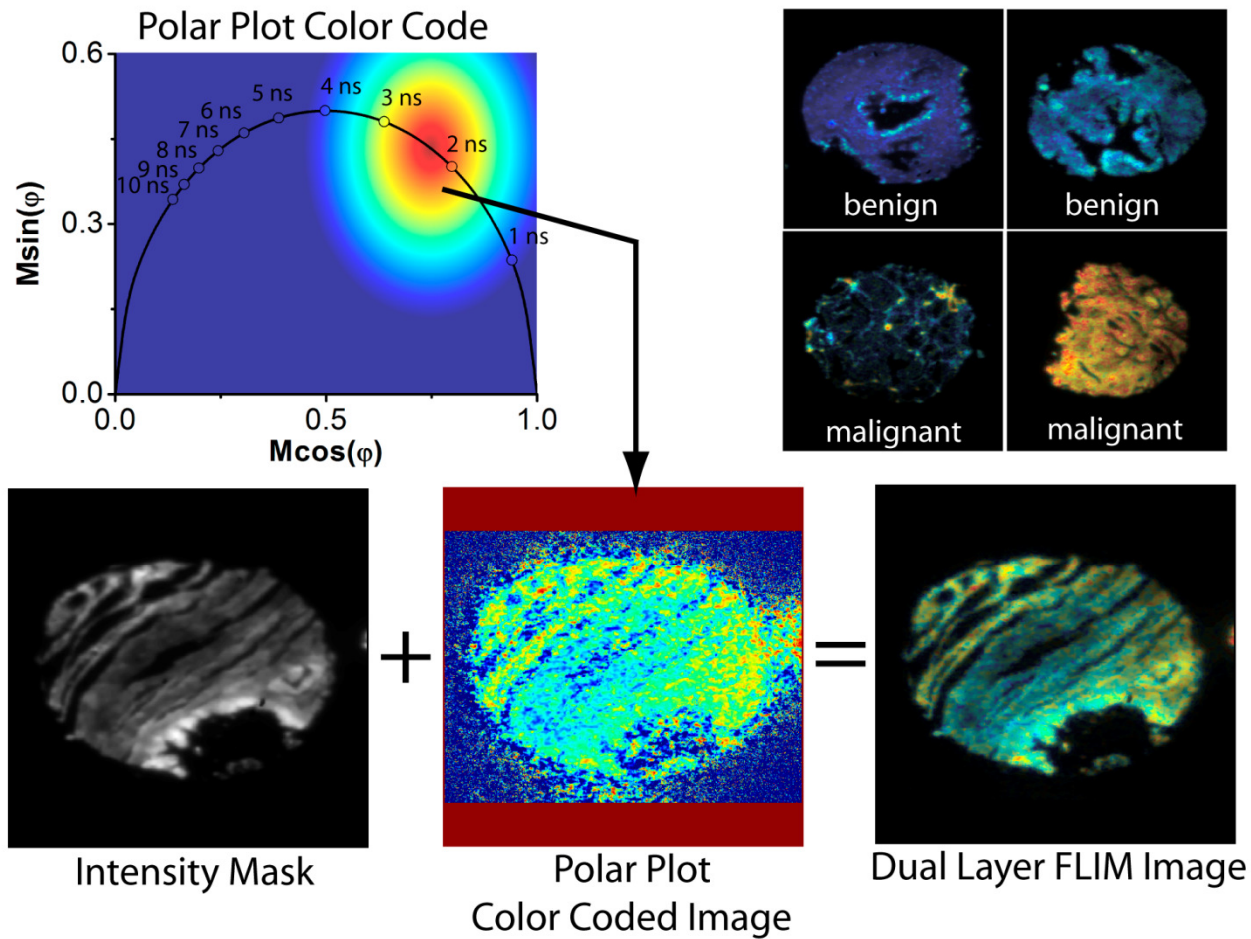
Line profile from an image of prostate tissue

Spatial frequency cuts (intervals) can be selected
Edges not smoothed

Rebecca Willett and Robert Nowak

Polar Plots

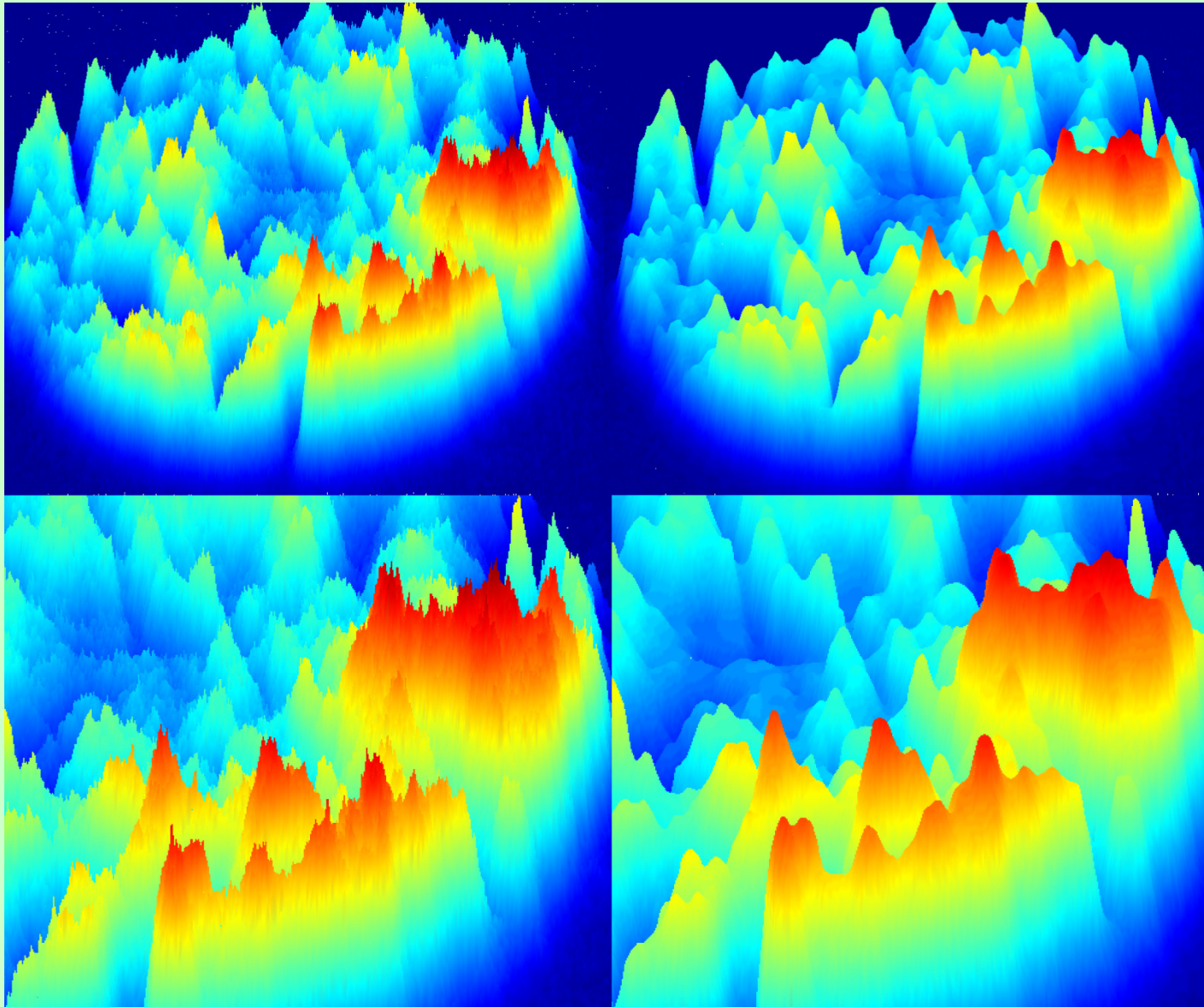


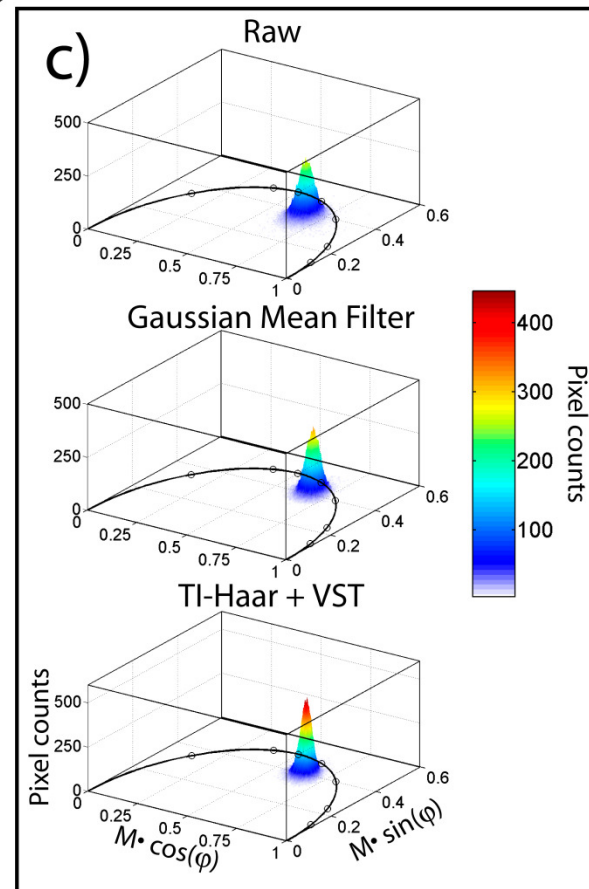
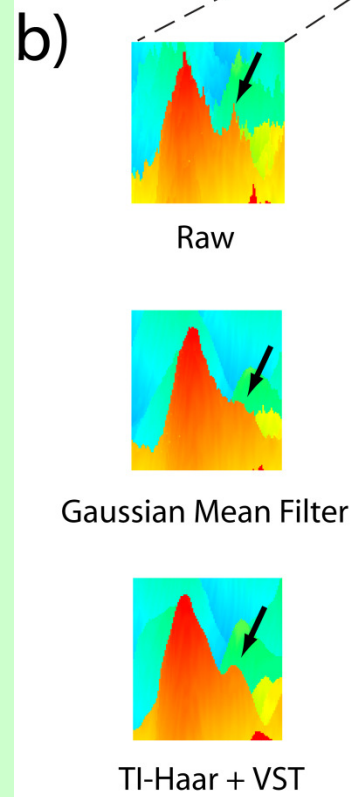
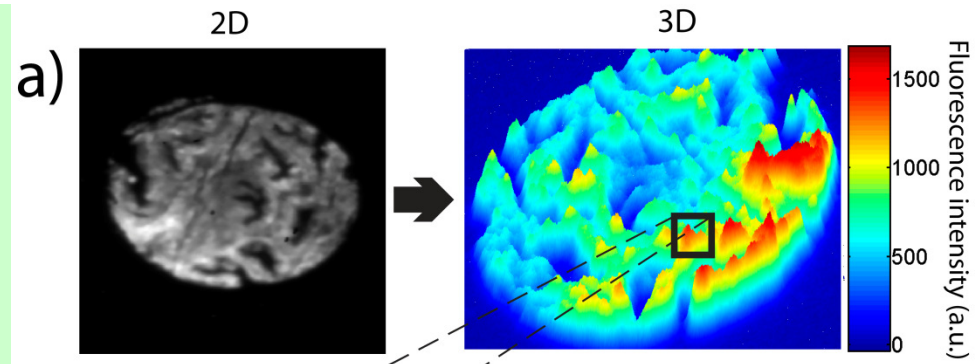


Combining morphological features and FLIM signals (wavelets)
and
Using denoising to assist in the overall analysis

Left: raw fluorescence
intensity images of a
prostate tissue core

Right: denoised images



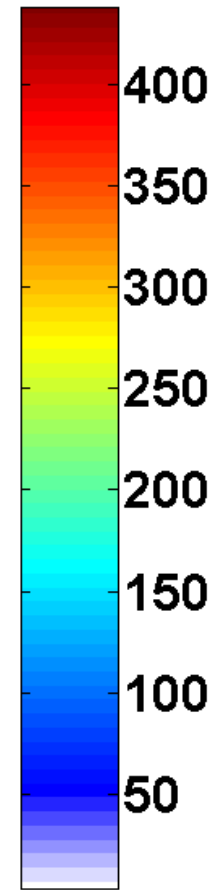
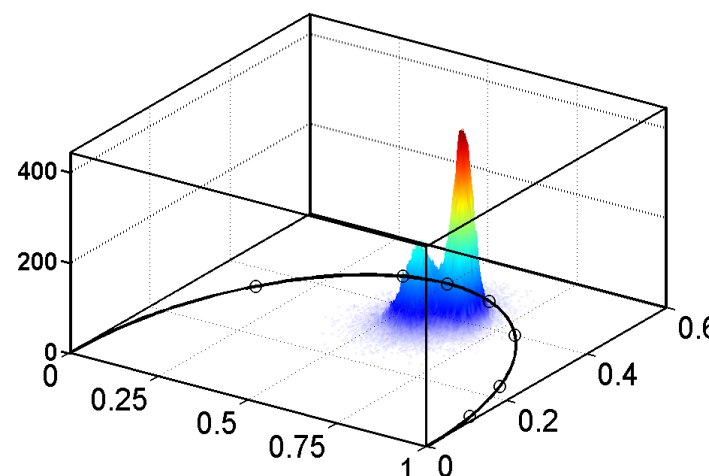
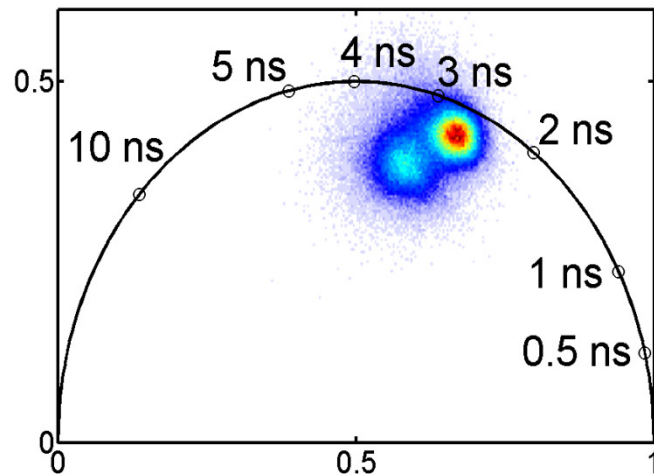
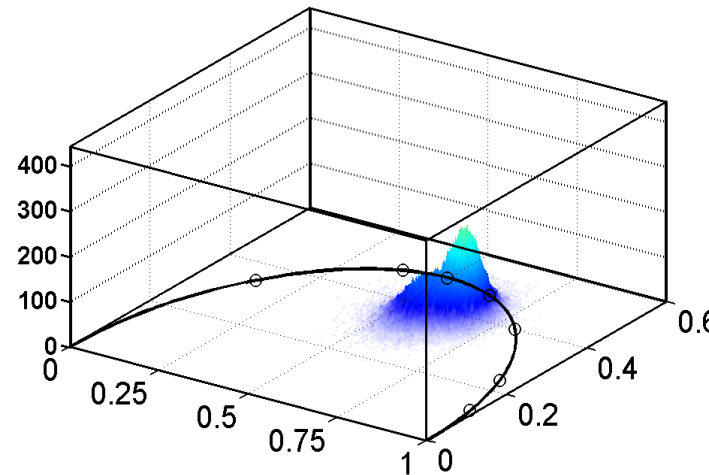
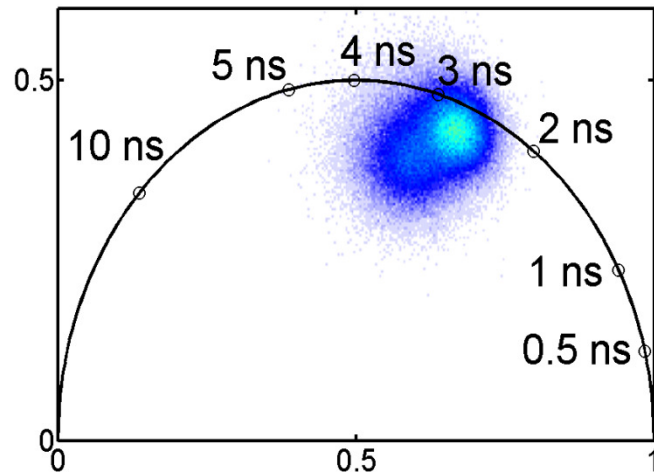


Polar plot histograms

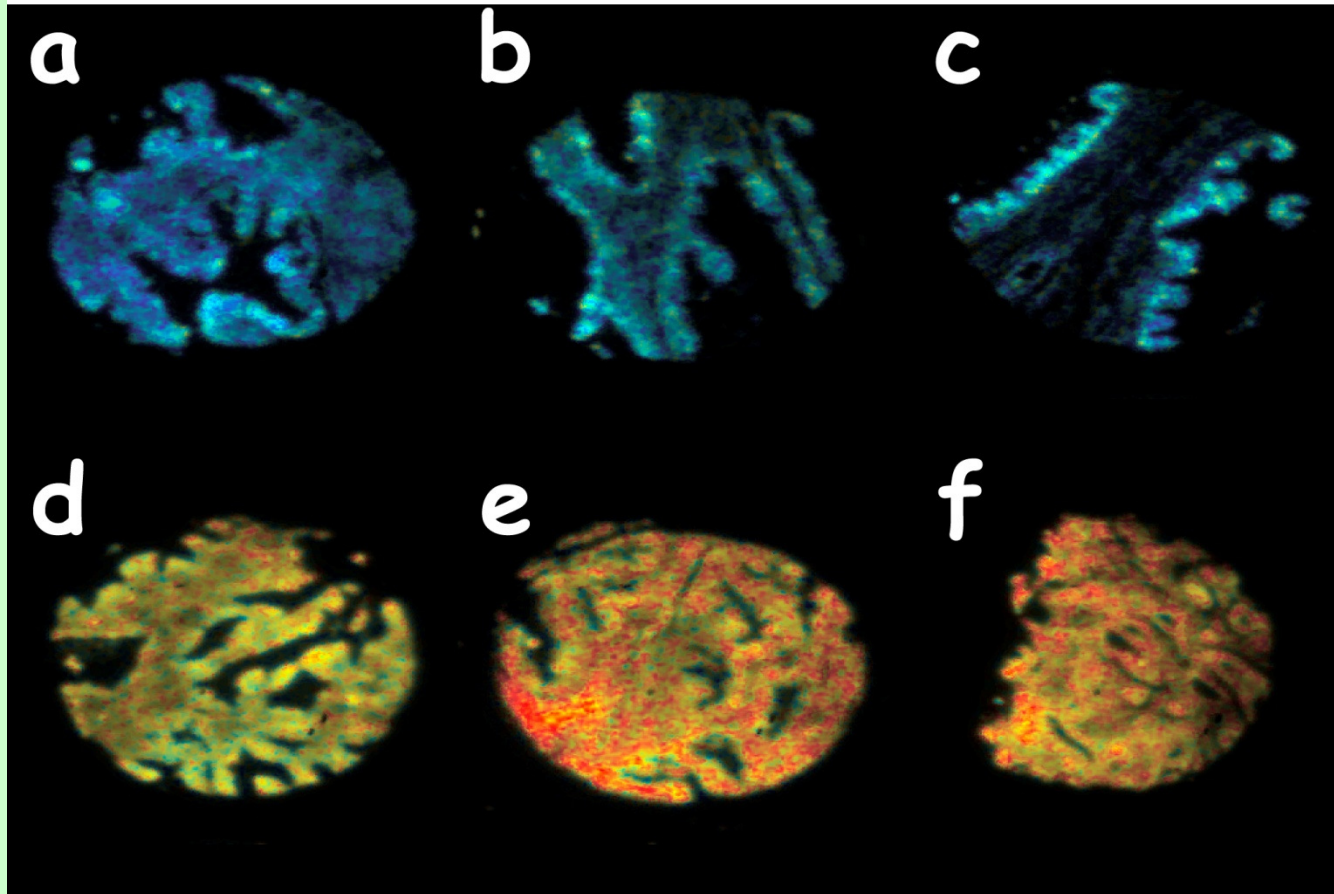
of entire FLIM images of a benign and a malignant prostate tissue core

Top: **before** denoising

Bottom: **after** denoising



Pixel Counts



Dual layer FLIM images (intensity image used to mask color coded lifetime image) - Color indicates the fluorescence lifetime distribution of each pixel. Blue indicates normal fluorescence lifetimes while red indicates a significant shift in fluorescence lifetimes from benign tissue. **a-c**) Benign tissue cores. **d**) Low-grade cancer. **e & f**) High-grade cancer. Note that the lifetime distributions can be complex (the fluorescence lifetimes reflect multiple species – i.e., free and enzyme-bound species). The color coding represents an overall shift in the relative amounts of each species and therefore accomplishes representation of complex data in an easily visible fashion.

SO.....

With FLIM
We need all the
help we can get.
Morphology helps!



Full Field FLI

Peter Schneider

Oliver Holub

Christoph Gohlke

Glen Redford

(polar plot ALA-PPIX)

+

Spinning disk, wavelets and denoising

Chittaton Buranachi, Bryan Spring, Rohit Bhargava

(dendritesALA-PPIX, prostate FLIM, redox sensor)

Yi-Chun Chen (Polar Plot, spectral FLIM, photosynthesis)

Photosynthesis:

Govindjee

Oliver Holub

Christoph Gohlke

Gregor Heiss

Shizue Matsubara

